



Computational Challenges in High Intensity and High Brightness Beam Physics

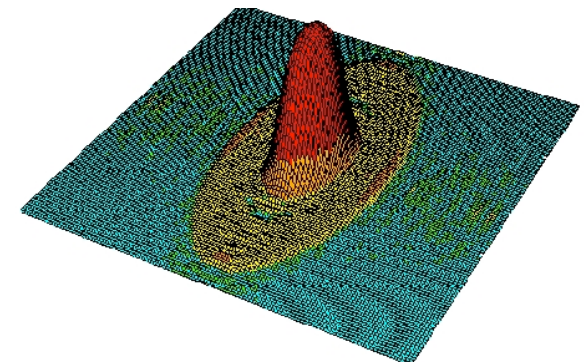
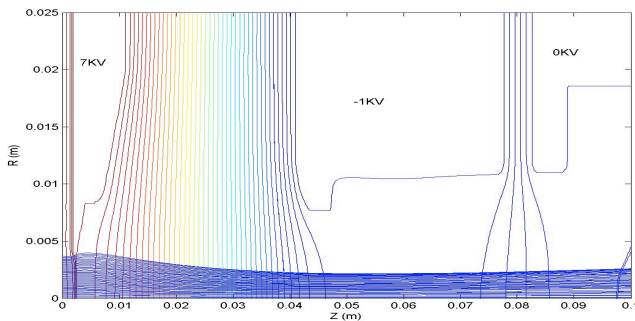
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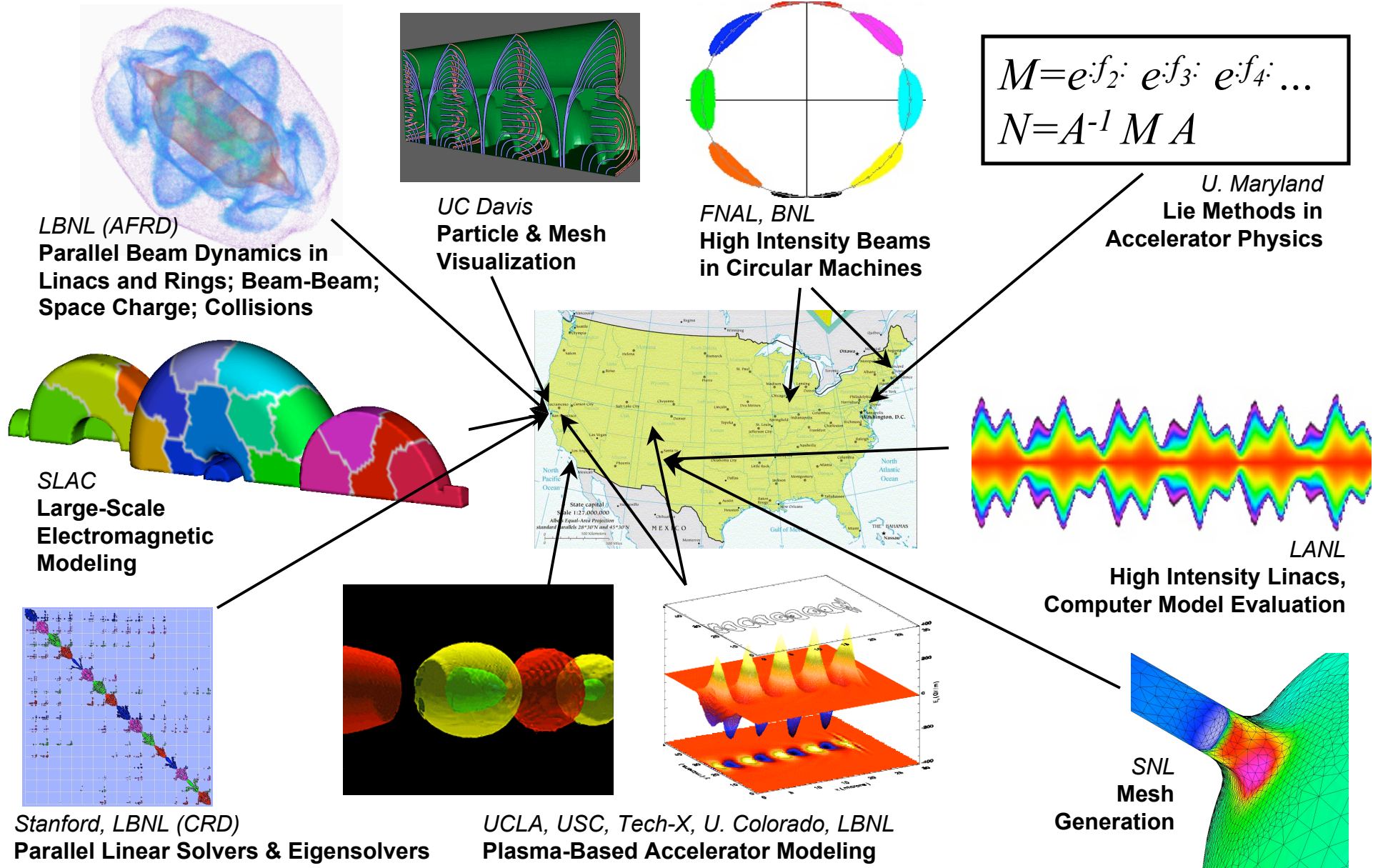




Acknowledgement

**Thanks to
Ji Qiang, LBNL**

SciDAC AST: A multidisciplinary, multi-institutional effort to produce comprehensive parallel accelerator design tools



SciDAC Applications



- **Beam dynamics codes developed under the SciDAC AST project have been applied to many projects. Some examples**
 - Tevatron, RHIC, PEP-II, LHC
 - NLC, ILC
 - FNAL Booster
 - RIA
 - SNS, LCLS
 - CERN SPS benchmark study
 - CERN SPL design
 - J-PARC commissioning

Computational beam dynamics: toward higher intensity, higher brightness, and greater precision



- **Modeling challenges for next-generation light sources:**
 - Photoinjector design and limits to brightness
 - Emittance preservation in the presence of **CSR**
 - Issues for energy recovery linacs
 - Diagnostics: design of ultrafast streak cameras
 - ...
 - Modeling challenges for high intensity rings (e.g. accumulators):**
 - Predicting halo formation, beam loss, and stability thresholds
- **Modeling challenges for other rings**
 - Predicting long-term effect of weak space charge + nonlinearities
- **Modeling challenges for colliders**
 - Maximizing luminosity, predicting lifetime due to b-b effect

Predicting radiation production/consequences



- Sub-ps duration bunches emit “long” wavelength (i.e. $\lambda \sim 100 \text{ nm} - 100 \mu\text{m}$) coherent radiation in bends (i.e. CSR) and in undulators (CUR, or coherent undulator radiation)
- This radiation can cause **microbunching** instabilities, increased energy spread, and emittance growth --- can strongly reduce the *desired* FEL instability in x-ray devices such as the LCLS
- Some CSR modeling capabilities rely on N^2 algorithms; leads to
 - numerically noisy simulations
 - limited spatial & temporal resolution
- Minimal CUR modeling capability when considering interaction with simultaneously growing FEL instability - -- vast longitudinal scale length difference.
Multiscale/multiresolution techniques needed.

Computational Challenges



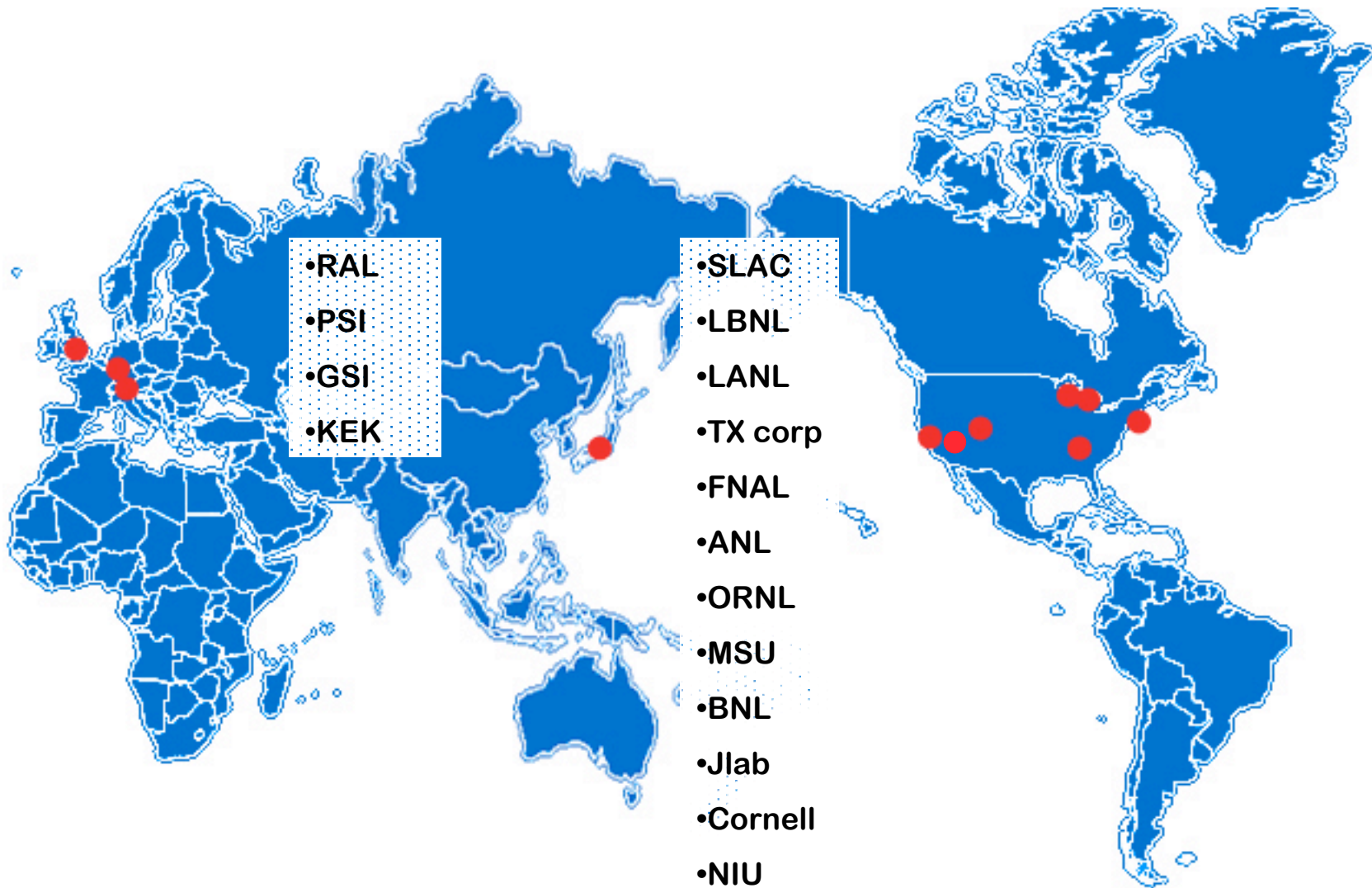
- **Multi-physics modeling:**
 - **Nonlinear optics, space charge, wakes, CSR, collisions**
- **Multi-scale modeling:**
 - **In space and time**
- **Getting the physics right. For example:**
 - **Tools for computing transfer maps from field data**
 - **Self-consistent modeling (when needed)**
- **Using codes to make predictions**
 - **Code benchmarking, code calibration, combining simulation and experiment to make predictions**
- **Visualizing/exploring massive data sets**
- **Building codes to make use of state-of-the-art platforms**
 - **Computers >10K processors here now, >100K soon**

IMPACT (Integrated-Map and Particle Accelerator Tracking) code (J. Qiang)



- A code suite (linac design code, 3D rms envelope code, 2 parallel PIC tracking codes)
- Recent enhancements
 - IMPACT-T: time-based (instead of z-based) version
 - Cathode emission model
 - Images from cathode
 - Poisson solver for high aspect ratio situations
 - Energy binning for large energy spread
 - Multi-charge state capability
- These enhancements benefit projects
 - Modeling rf photoinjectors
 - Modeling beams w/ large energy spread, as for electron bunches emerging from plasma-based accelerators
 - Multi-charge state for RIA

Impact User-Map



Photoinjector Modeling



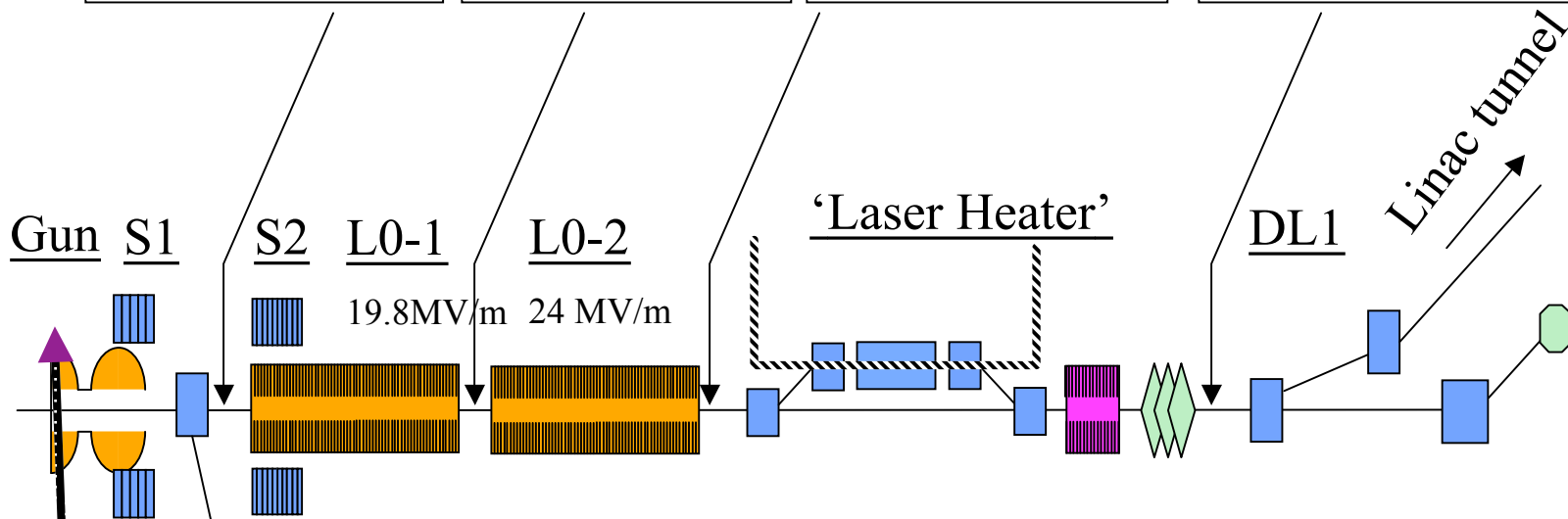
Several groups are using IMPACT-T (or have expressed an interest) for photoinjector studies

- Fermilab RF photoinjector
- LCLS
- Fermi@Elettra project
- Cornell ERL
- ALS streak camera
- BNL electron cooling project
- ANL advanced accelerator project
- JLab

A Schematic Plot of LCLS Injector Layout



6 MeV $\epsilon = 1.6 \mu\text{m}$ $\sigma_{\delta,un.} = 3\text{keV}$	63 MeV $\epsilon = 1.08 \mu\text{m}$ $\sigma_{\delta,un.} = 3\text{keV}$	135 MeV $\epsilon = 1.07 \mu\text{m}$ $\sigma_{\delta,un.} = 3\text{keV}$	135 MeV $\epsilon = 1.07 \mu\text{m}$ $\sigma_{\delta,un.} = 40\text{keV}$
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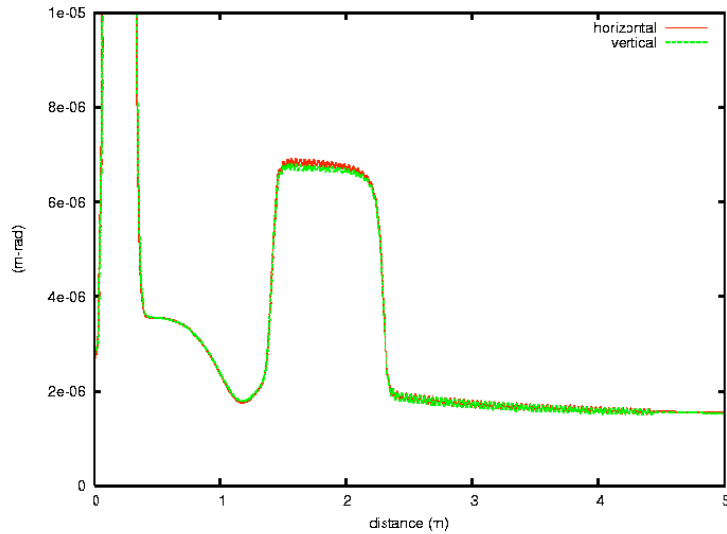
UV Laser 200 μJ ,

$\lambda = 255 \text{ nm}$, 10ps, $r = 1.2 \text{ mm}$

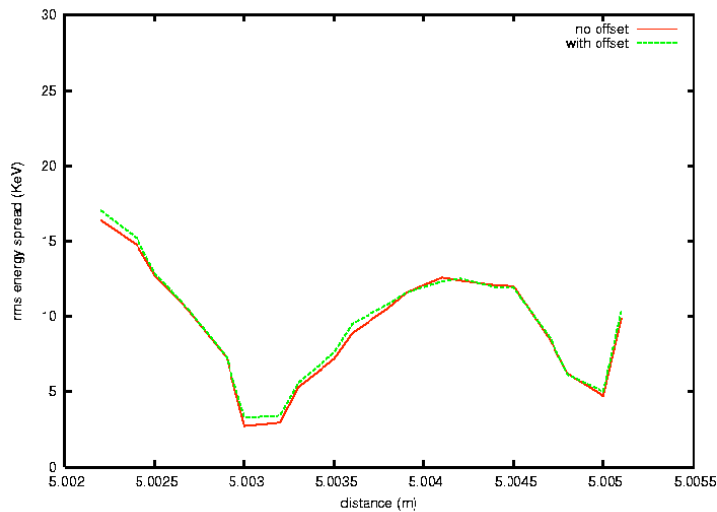
Peak Current	100 A
Charge	1 nC
$\epsilon_{\text{projected}}, \epsilon_{\text{slice}}$	$\leq 1.2, 1 \mu\text{m}\cdot\text{rad}$
Repetition rate	120 Hz

LCLS modeling using IMPACT-T

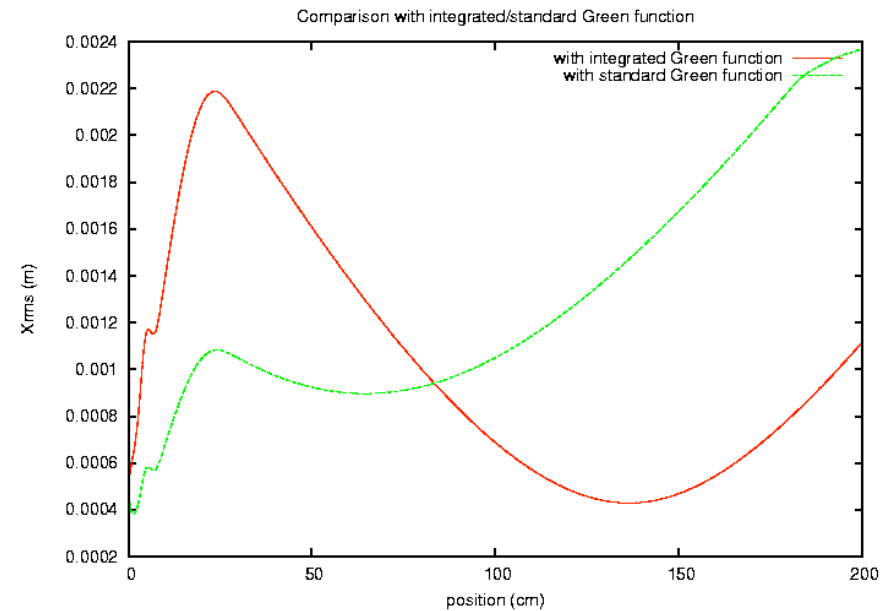
J. Qiang/LBNL and C. Limborg/SLAC



LCLS injector transverse emittance vs distance



LCLS slice rms energy spread after 1st TW linac (w/ and w/out initial 100 μ offset)

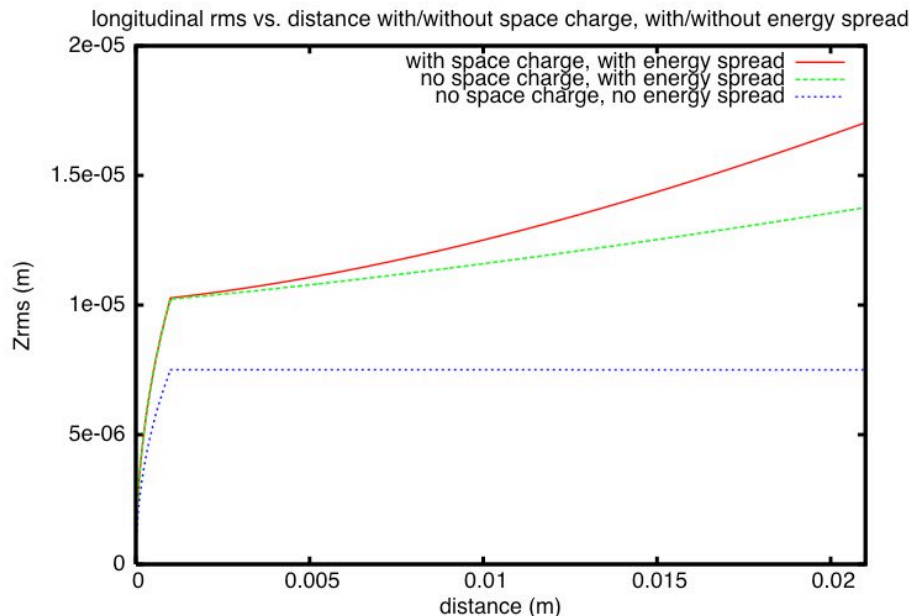


LCLS simulation: Large effect observed when using Integrated Green function compared with standard Green function

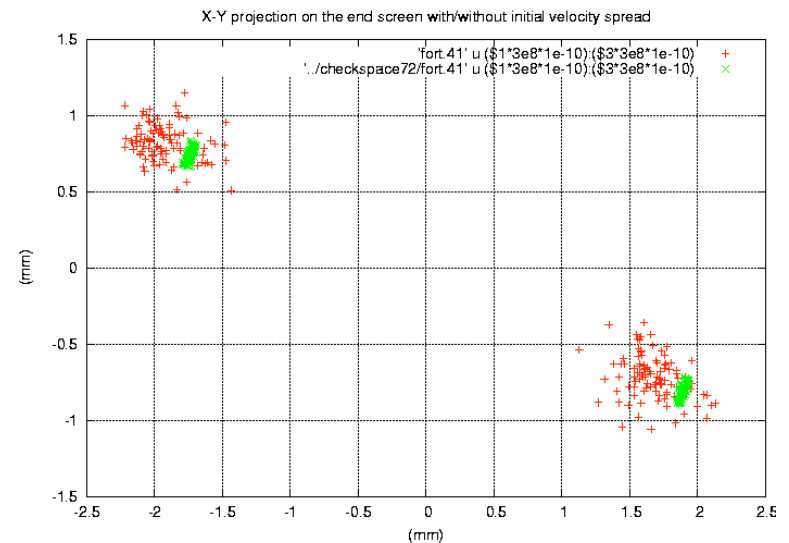
Design of ultrafast streak cameras



- Ultrafast streak cameras being designed for LCLS and ALS
 - Simulating LCLS version challenging: space-charge effects
 - Simulating ALS version challenging: small # of electrons (N^2 code)
- Both currently being modeling using IMPACT-T



Plot of rms bunch length vs distance in the proposed LCLS streak camera showing the effect of space charge on the bunch length



X-Y projection on the end screen w/ and w/out energy spread in the ALS streak camera (proposed LBNL LDRD)

Algorithmic Advances



- **Integrated Green function**
 - Solves the long-standing problem that has plagued certain codes (e.g. PARMELA, original IMPACT) when grid aspect ratios become large
- **Shifted Green function**
 - Originally developed to model long-range beam-beam interactions
 - Same algorithm used to treat cathode image effects
- **Wavelet-based methods**
- **Multi-level gridding, AMR**

Integrated Green Function (IGF) addresses a Critical Issue: high aspect ratios



- Poisson solvers used in static electric and gravitational particle-in-cell simulations **often fail** when the grid aspect ratio $\gg 1$
- **Some important problems involve extreme aspect ratios:**
 - Long beams in rf accelerators; pancake beams
 - Beams in induction linacs: $L \sim 10$ s of meters; $R \sim \text{cm}$
 - Galaxies
- Standard grid-based approaches involve using a very large # of grid points in the long dimension, leading to prohibitively long run times
 - As a result, it is *extremely* difficult model high aspect ratio systems accurately using standard grid-based approaches

IGF recognizes that certain physical quantities may vary on vastly different scales



- The Green function, G , and source density, ρ , may change over different scales
- G is known a priori; ρ is not

We should use all the information available regarding G so that the numerical solution is only limited by our approximate knowledge of ρ

- Example: 2D Poisson equation in free space

$$\phi(x, y) = \int G(x - x', y - y') \rho(x', y') dx' dy'$$

Standard Approach (Hockney and Eastwood)



$$\phi_{i,j} = \sum G_{i-i',j-j'} \rho_{i',j'}$$

- This approach is equivalent to using the trapezoidal rule (modulo treatment of boundary terms) to approximate the convolution integral
- This approach makes use of only partial knowledge of G
- The error depends on how rapidly the integrand, ρG , varies over an elemental volume
 - If ρ changes slowly we might try to use a large grid spacing; but this can introduce huge errors due to the change in G over a grid length

IGF Algorithm



- Assume the charge density, ρ , varies in a prescribed way in each cell
- Use the analytic form of the Green function to **perform the convolution integral exactly** in each cell, then sum over cells
- Example: linear basis functions to approximate ρ in a cell:

$$\begin{aligned} \phi(x_i, y_j) = & \sum_{i', j'} \rho_{i', j'} \int_0^{h_x} dx' \int_0^{h_x} dy' (h_x - x')(h_y - y') G(x_i - x_{i'} - x', y_j - y_{j'} - y') + \\ & \sum_{i', j'} \rho_{i+1, j} \int_0^{h_x} dx' \int_0^{h_x} dy' x' (h_y - y') G(x_i - x_{i'} - x', y_j - y_{j'} - y') + \\ & \sum_{i', j'} \rho_{i, j+1} \int_0^{h_x} dx' \int_0^{h_x} dy' (h_x - x') y' G(x_i - x_{i'} - x', y_j - y_{j'} - y') + \\ & \sum_{i', j'} \rho_{i+1, j+1} \int_0^{h_x} dx' \int_0^{h_x} dy' x' y' G(x_i - x_{i'} - x', y_j - y_{j'} - y') \end{aligned}$$

- Shifting the indices results in a single convolution involving an integrated effective Green function: $\phi_{i,j} = \sum G_{i-i', j-j'}^{eff} \rho_{i', j'}$

Cost and Accuracy of IGF; Improvement over Standard Approach



- **Cost:** IGF elemental integrals can be done analytically; formulas are very lengthy
 - Requires more FLOPS than simply using G_{ij} but...
 - In situations where the grid is fixed, this only needs to be done once at the start of a run. Amortized over many time steps, this does not significantly impact run time.
- **Accuracy:** Method works as long as the elemental integrals are computed accurately and as long as the grid and # of macroparticles are sufficient to resolve variation in ρ
- IGF maintains accuracy even for extreme aspect ratios (>1000:1)

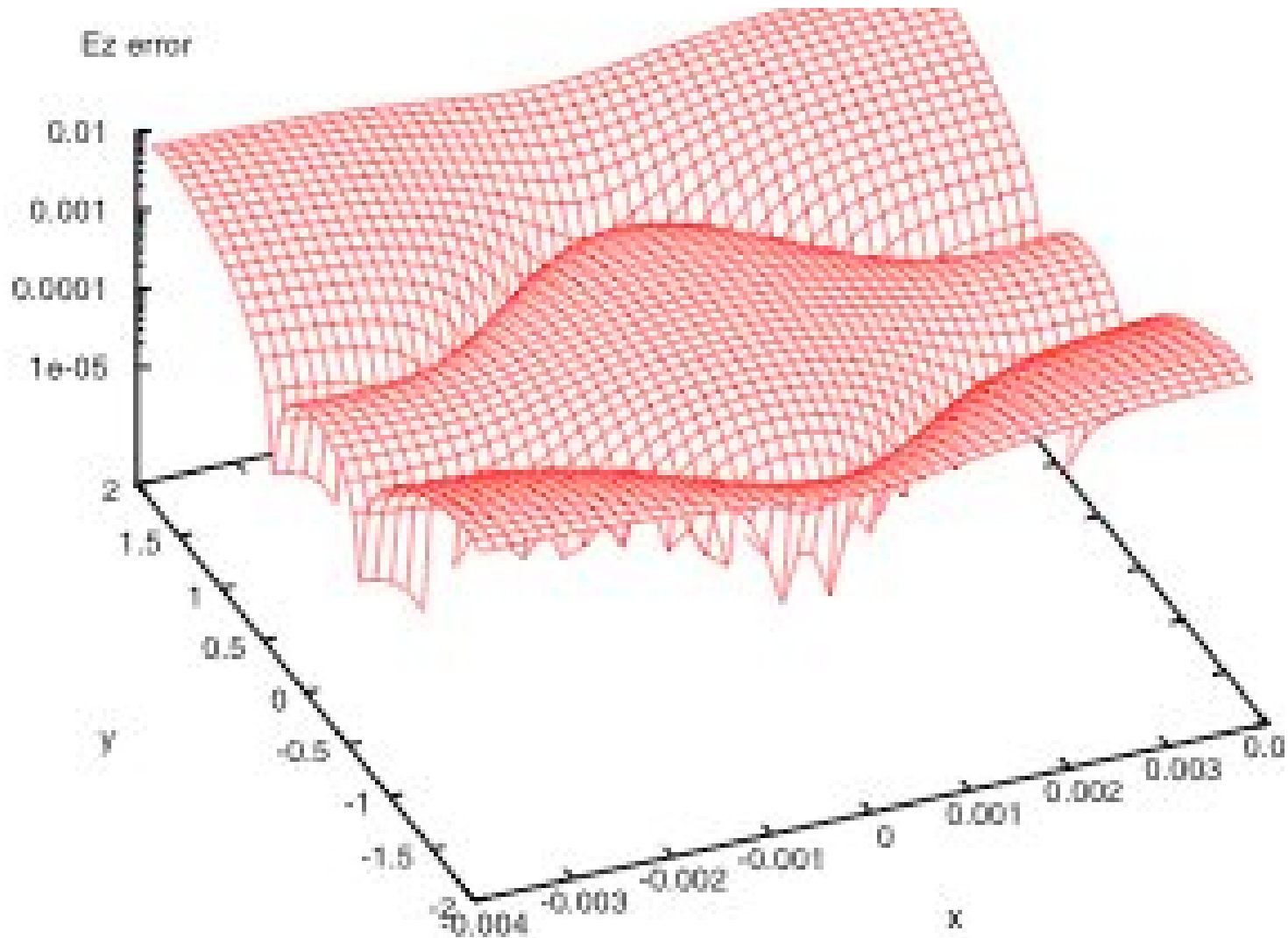
As a result, IGF performs orders of magnitude better than the standard convolution algorithm for realistic problems involving large aspect ratios

Example: 2D gaussian ellipse



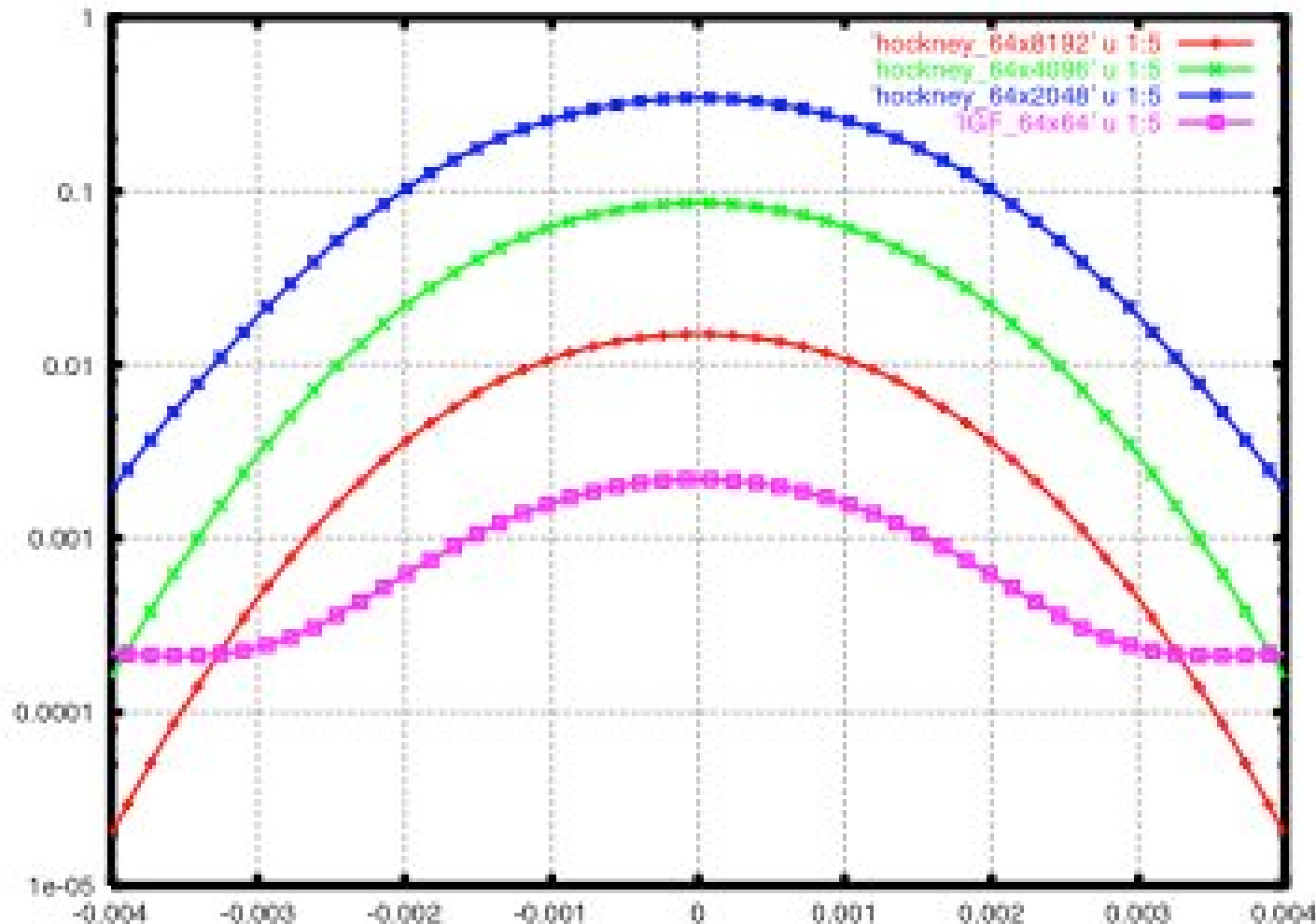
- **Aspect ratio is 1:500 -- $x_{\max}=0.002$, $y_{\max}=1$**
- **Calculation of fields using (1) standard Hockney algorithm and (2) IGF approach**
 - In both cases, performed convolutions for the fields directly (rather than calculating the potential and using finite differences to obtain fields)**
- **Calculation performed using a mesh of size**
 - Hockney: 64x64, 64x128, 64x256, ..., 64x16384**
 - IGF: 64x64**

IGF field error



Electric field error using IGF is below 1% using a 64x64 grid.

Comparison of IGF vs standard Hockney approach



Simulation of a high-aspect ratio bunch using an integrated Green Function (IGF) and a conventional algorithm (Hockney). IGF on a 64x64 grid (purple) is more accurate than a standard calculation using 64x2048 (blue), 64x4096 (green), and 64x8192 (red).

Conclusion: 2D Gaussian example



- **For this test problem, the standard Hockney algorithm would require ~500 times more computational effort to achieve the same worst-case accuracy as a simulation using the IGF approach.**
- **IGF works whether the aspect ratio is large, small, or near unity, i.e. it is generally applicable.**
- **Verified in 2D and 3D**
 - 3D might require quad precision

Extension of IGF to Beams in Pipes



- IGF is especially useful when applied to beams in pipes, since the Green function falls off exponentially in z , though $\rho(z)$ may change slowly over meters
- Due to shielding in beampipe, sum can be truncated in the “long” direction:

$$\phi_{i,j} = \sum_{i'=1}^{N_x} \sum_{j'=j}^{j \pm j_{cutoff}} G_{i-i', j-j'}^{eff} \rho_{i',j'}$$

- If grid length in z is \gg pipe radius, can truncate at nearest neighbors:

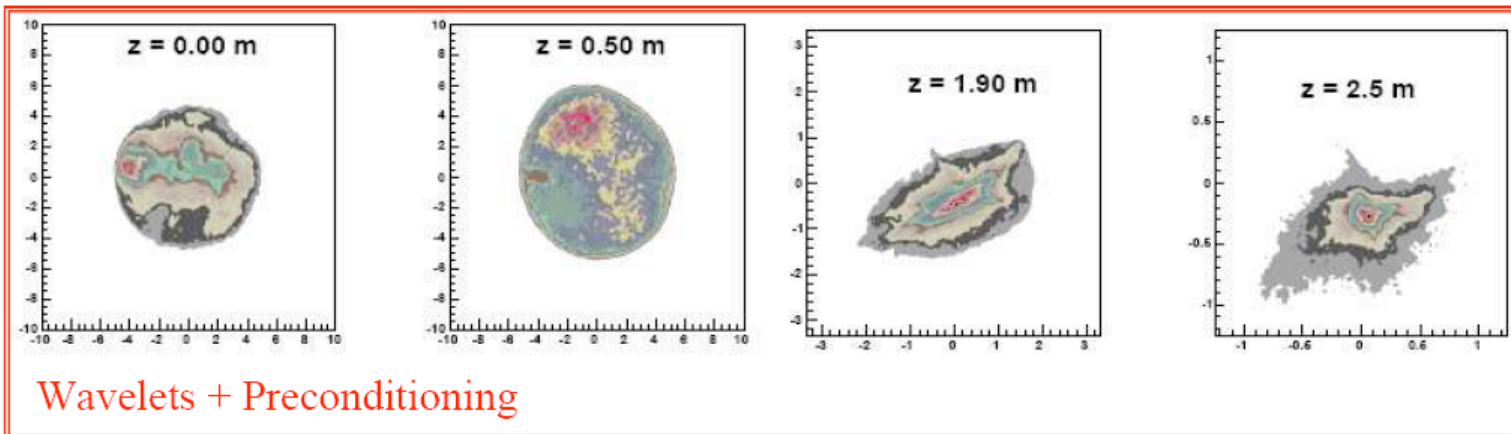
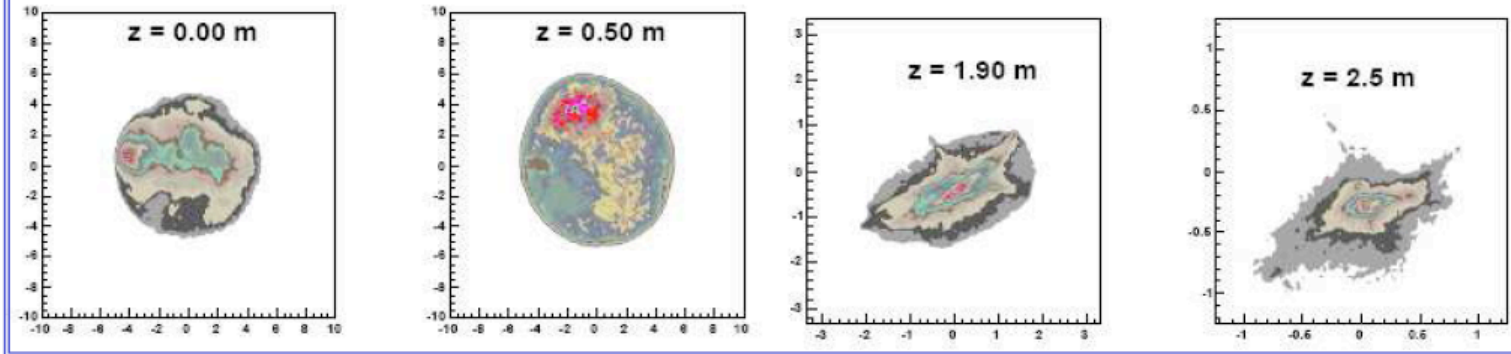
$$\phi_{i,j} = \sum_{i'=1}^{N_x} (G_{i-i', j-1}^{eff} \rho_{i',j-1} + G_{i-i', j}^{eff} \rho_{i',j} + G_{i-i', j+1}^{eff} \rho_{i',j+1})$$

- For a rectangular pipe, can rewrite Green function as a sum of convolutions and correlations; can still use FFT-based approach to sum over elements
- Applicability to circular pipes is still an open problem

Toward multiscale/multiresolution modeling: high level of spatial resolution maintained by wavelet-based solver



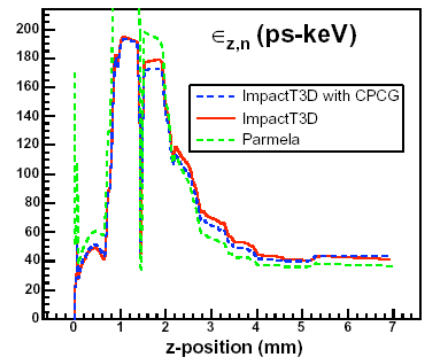
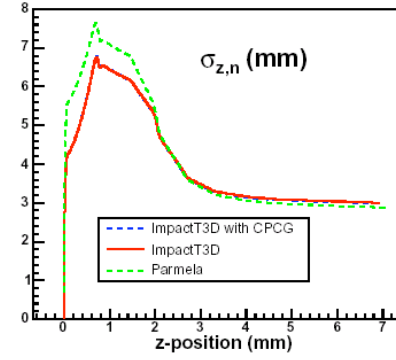
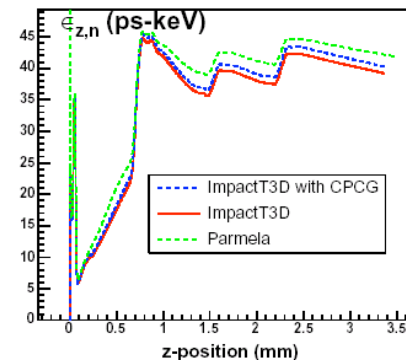
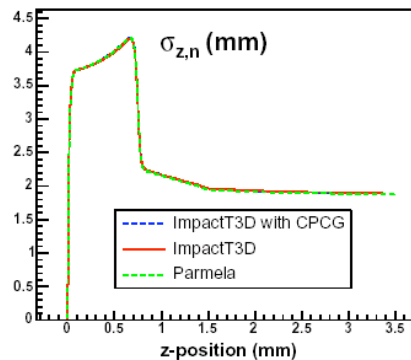
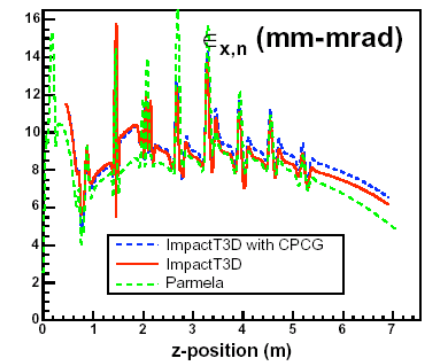
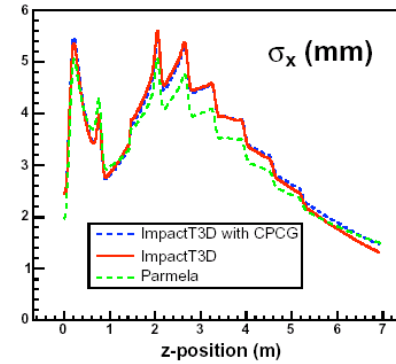
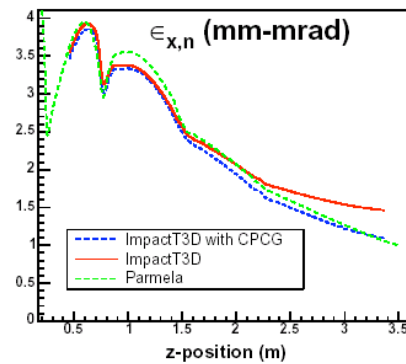
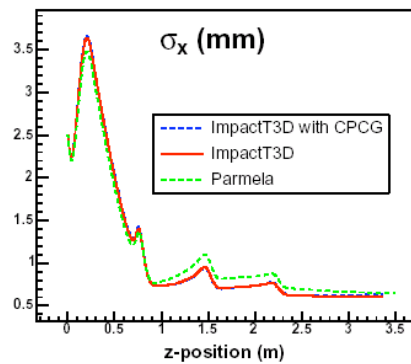
Green Function + Fast Fourier Transform



Comparison between IMPACT-T (upper plots), IMPACT-T w/ wavelet-based Poisson solver (lower plots) for transverse distributions at different z-locations in the AES/Jlab low-charge photoinjector. Wavelet solver is already as fast as standard solver in IMPACT-T; significant speedup will follow implementation of wavelet thresholding. Also investigating more effective preconditioner.

Simulation of AES/Jlab photoinjector

I. Pogorelov/LBNL and B. Terzek/Northern Illinois Univ.

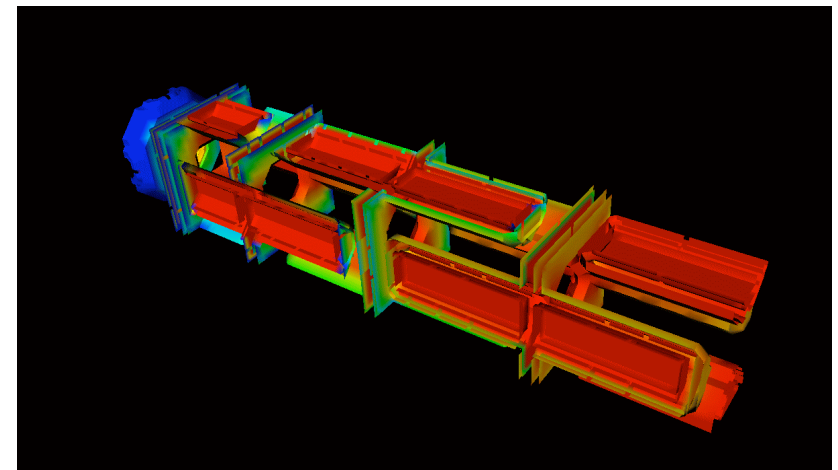
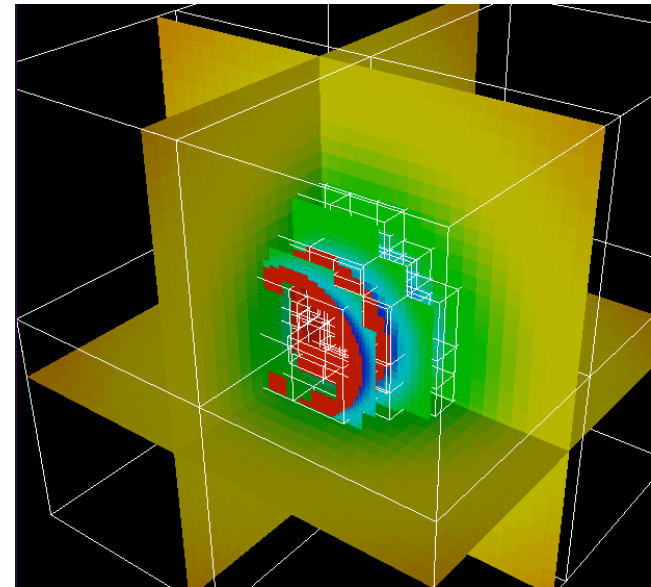


Comparison of IMPACT-T, IMPACT-T w/ iterative wavelet-based Poisson solver, and PARMELA for the AES/Jlab low-charge photoinjector (left 4 plots) and high-charge photoinjector (right 4 plots)

Particle-In-Cell AMR: Highly challenging!



- One approach uses capability originally developed for the combustion community
 - Chombo framework for block-structured AMR
- Issues/Goals
 - Eliminating numerical self-forces in nonuniform grids
 - Fast infinite-domain boundary conditions.
 - Flexible specification of interaction between grid and particle data.
 - Accurate representation of complex geometries.



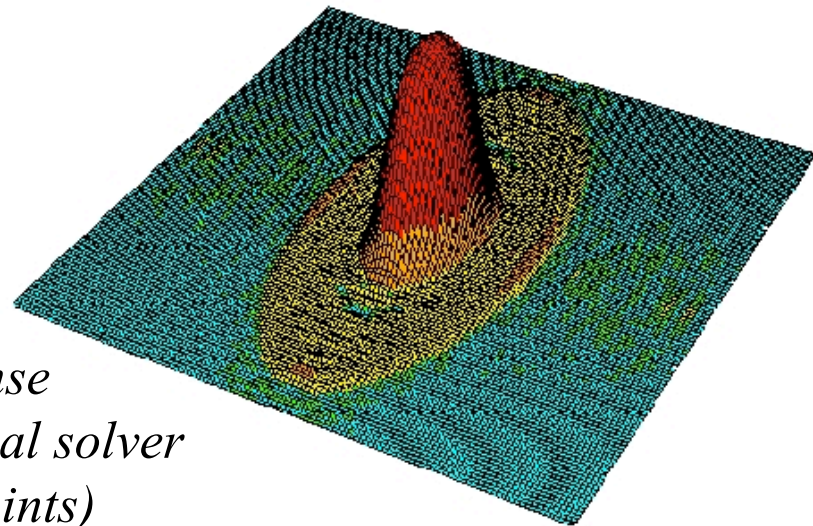
Direct Solvers: A future alternative to particle-based methods?



$$\frac{\partial f}{\partial t} + (\mathbf{p} \cdot \partial_{\mathbf{q}}) f - (\nabla \Phi \cdot \partial_{\mathbf{p}}) f = 0. \quad \text{Vlasov equation}$$

$$f(\mathbf{q}, \mathbf{p}, t) = \exp\left(-\frac{t}{2} \mathbf{p} \cdot \partial_{\mathbf{q}}\right) \exp(t \nabla \Phi \cdot \partial_{\mathbf{p}}) \exp\left(-\frac{t}{2} \mathbf{p} \cdot \partial_{\mathbf{q}}\right) f(\mathbf{q}, \mathbf{p}, 0) + O(t^3)$$

Split-operator stepping algorithm

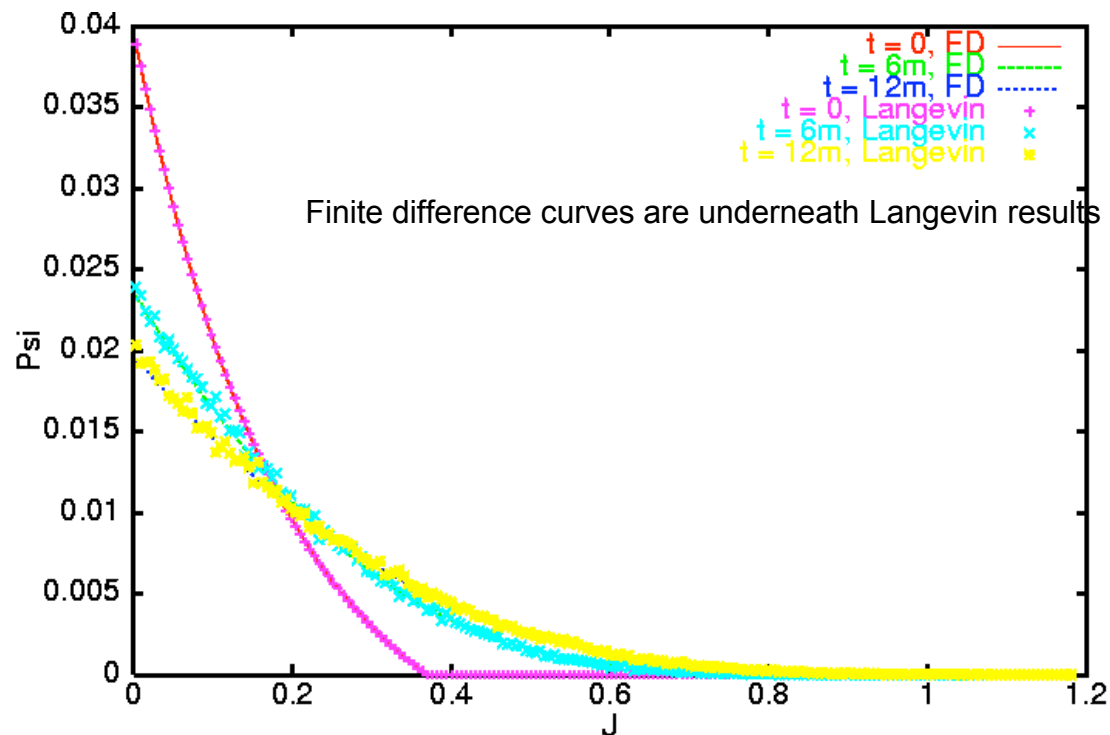


Density plot from a 2D simulation of an intense beam in a quadrupole channel using a spectral solver with a 128^4 phase-space grid (268M mesh points)

Modeling Additional Effects: Collisional Effects



- In collaboration w/ BNL (J. Wei), have performed comparisons of the self-consistent Langevin approach with a finite difference model
 - Langevin results and FD results are in excellent agreement, but Langevin does not have instability problem present in FD method

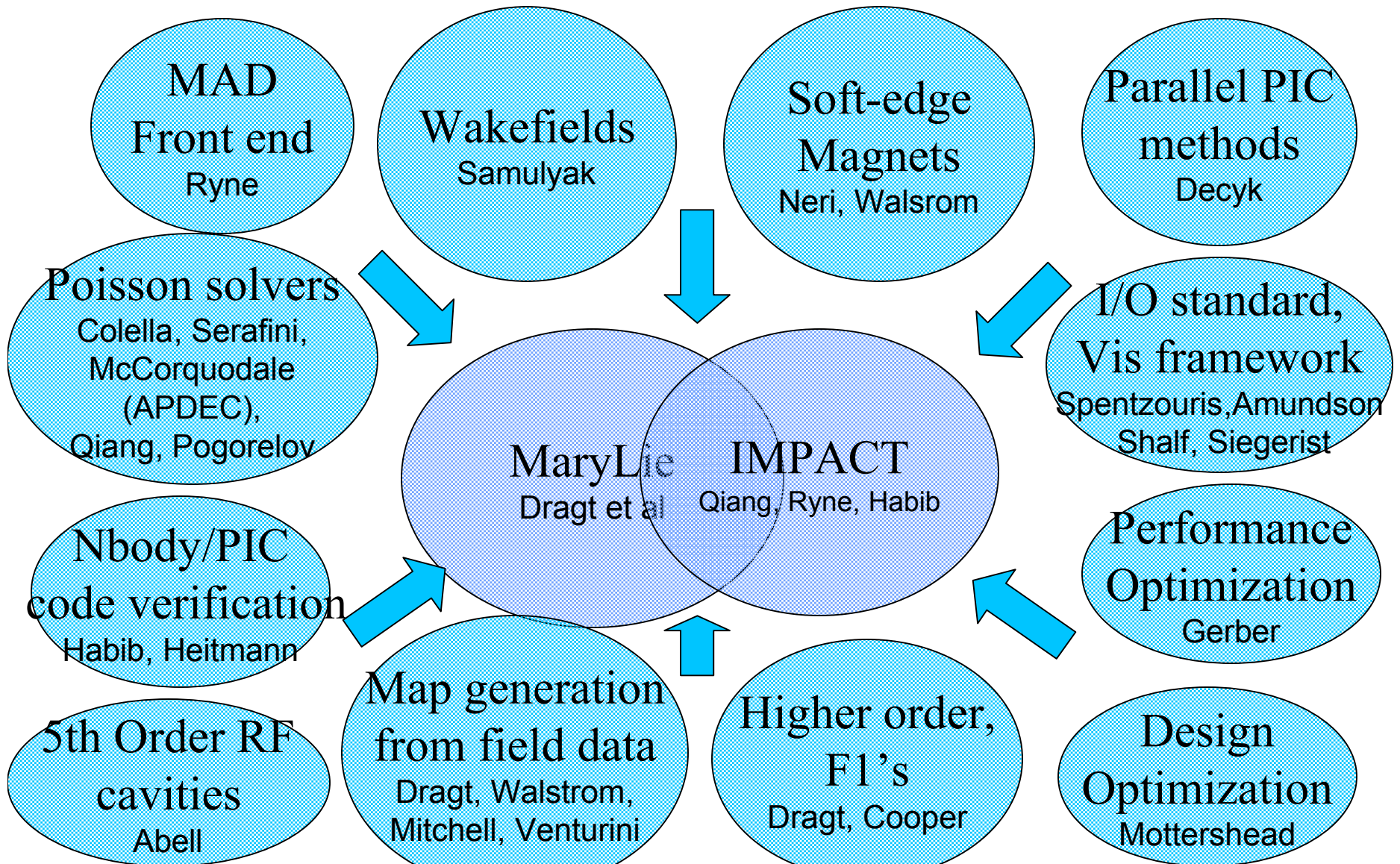


Code Componentization and Reuse



- **SciDAC model emphasizes large multidisciplinary teams; code componentization and reuse; involvement of applications scientists, computer scientists, mathematicians, visualization experts...**
- **Examples:**
 - Incorporation of Poisson solvers (P. Colella (APDEC), Chombo)**
 - Incorporation of PIC code parallelization strategies (V. Decyk, UPIC)**
 - Incorporation of parallel PIC data handling & viz (PARTVIEW, H5PART; J. Shalf, C. Siegerist, A. Adelman)**

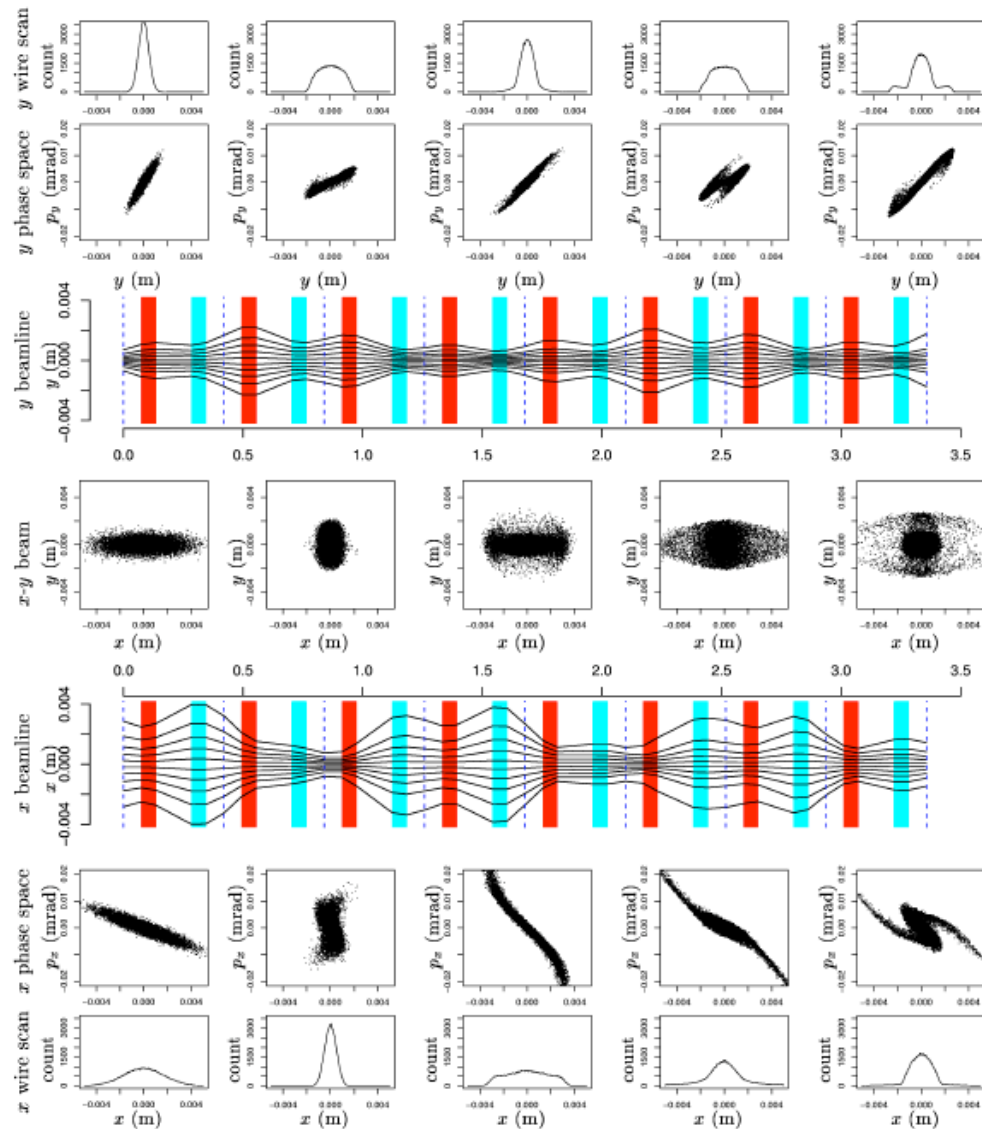
Code dev. involves large multidisciplinary teams. Example: MaryLie/IMPACT beam dynamics code



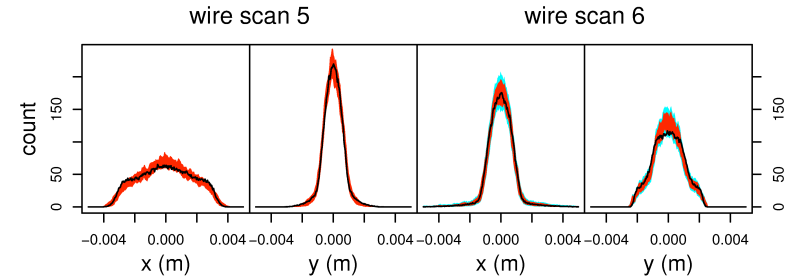


Predictability; Uncertainty Quantification; Code Verification and Validation

Statistical Methods for Calibration & Forecasting (D. Higdon et al, LANL)



- **Determining initial phase space distribution from 1D wire scan data.**



Simulation of a high intensity proton beam through a series of quadrupole magnets. Statistical techniques were used to combine 1D profile monitor data with simulations to infer the 4D beam distribution. The figure shows the 90% intervals for the predicted profile at scanner #6 (shaded regions), and, for comparison, the observed data (black line). Only data from the odd numbered scanners were used to make the prediction.

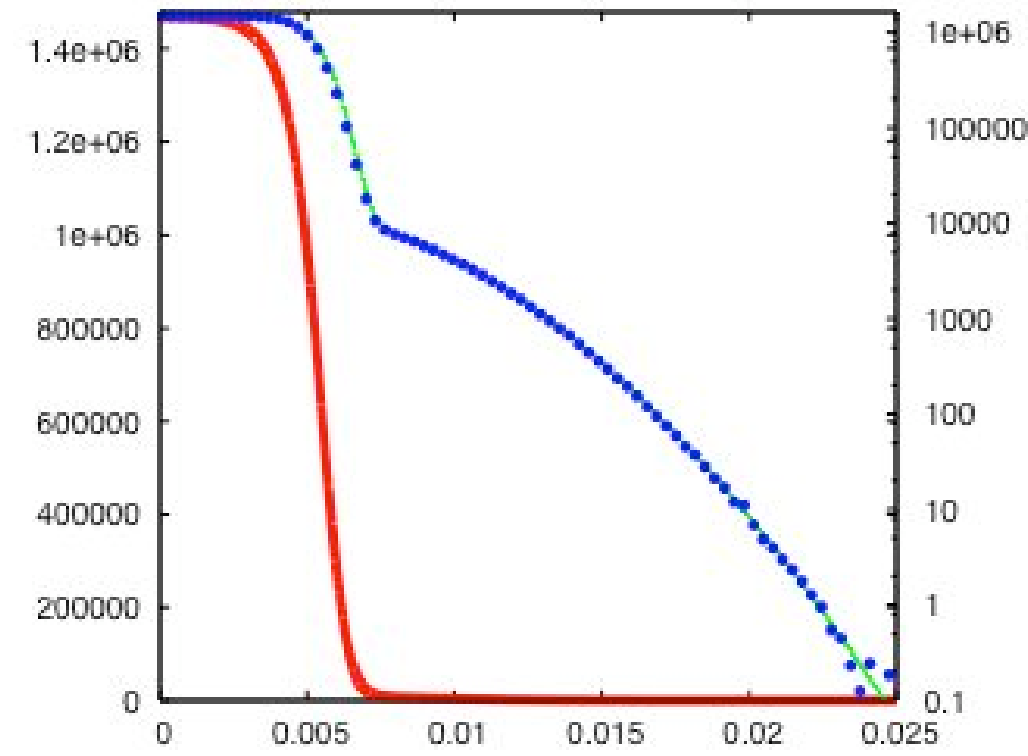
Code Test Suite

(dev in collab w/ A. Adelman, J. Amundson, P. Spentzouris)



- **KV beam in a FODO channel**
- **Free expansion of a cold, uniform density bunch**
- **Cold beam in a FODO channel with RF cavities**
- **Thermal beam in a constant focusing channel**
- **Bi-thermal beam in a constant focusing channel**

Bithermal distribution: A self-consistent 3D beam w/ a pronounced halo



ML/I simulation of bithermal distribution: 95% charge in core, 5% charge in halo. Note 6 order of magnitude resolution in simulated wirescan.

Question: Which accelerator physicist said this about N-body Simulations?



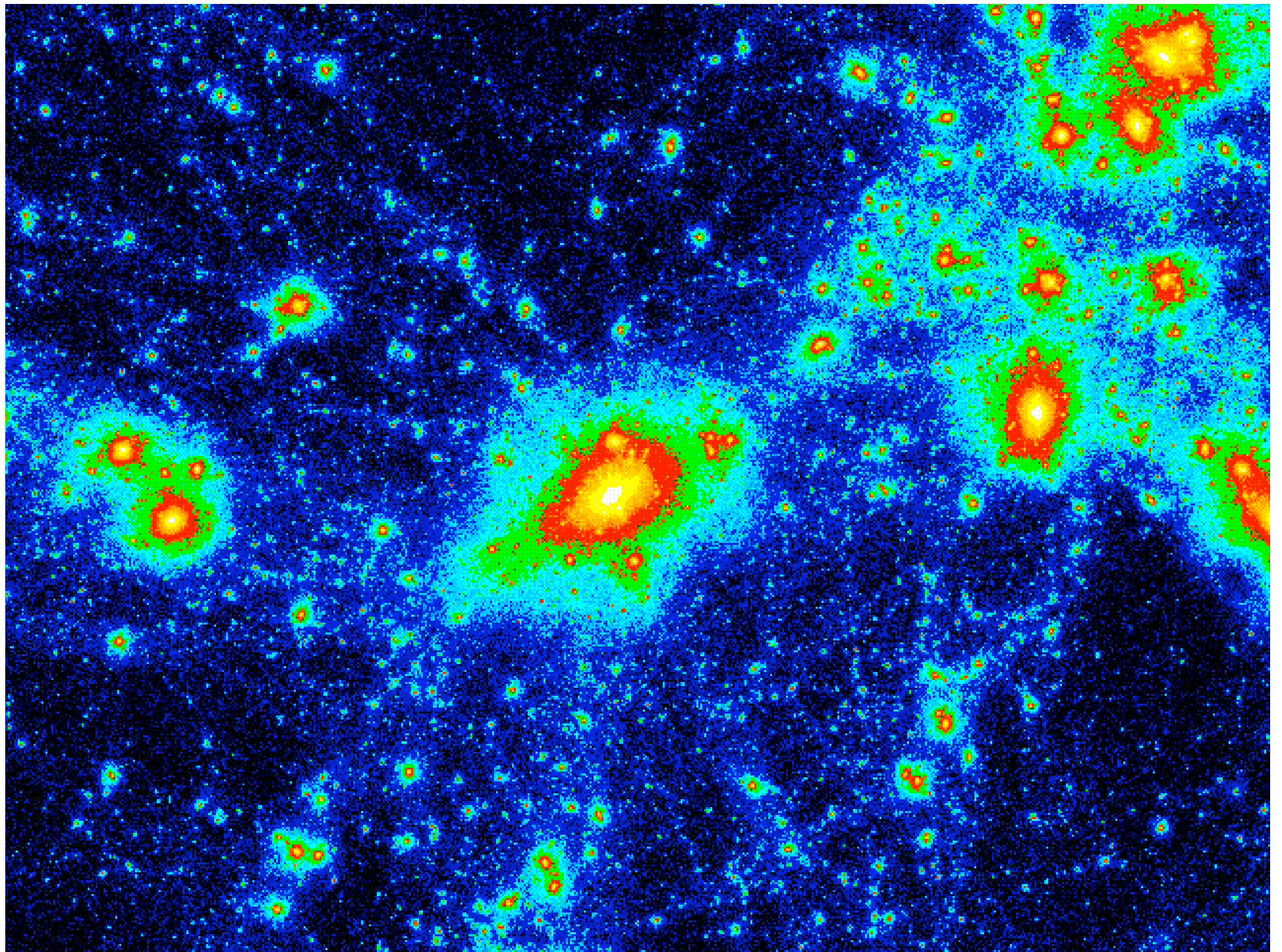
- **Task: solve collisionless Vlasov-Poisson equation; problems: 6-dim partial differential equation; because of nonlinearity ever smaller substructure is generated, instabilities occur as structure is generated on sub-resolution scales**
- **N-body approach: sample phase space distribution with tracer particles and evolve them by computing inter-particle forces -- N^2 problem**
- **further approximation techniques reduce force calculation to $\sim O(N)$ or $\sim O(N \log N)$; grid-based or multipole expansions or combination of both**
- **Klypin/Shandarin (1983): 32,768 particles; now: 1 billion !!**

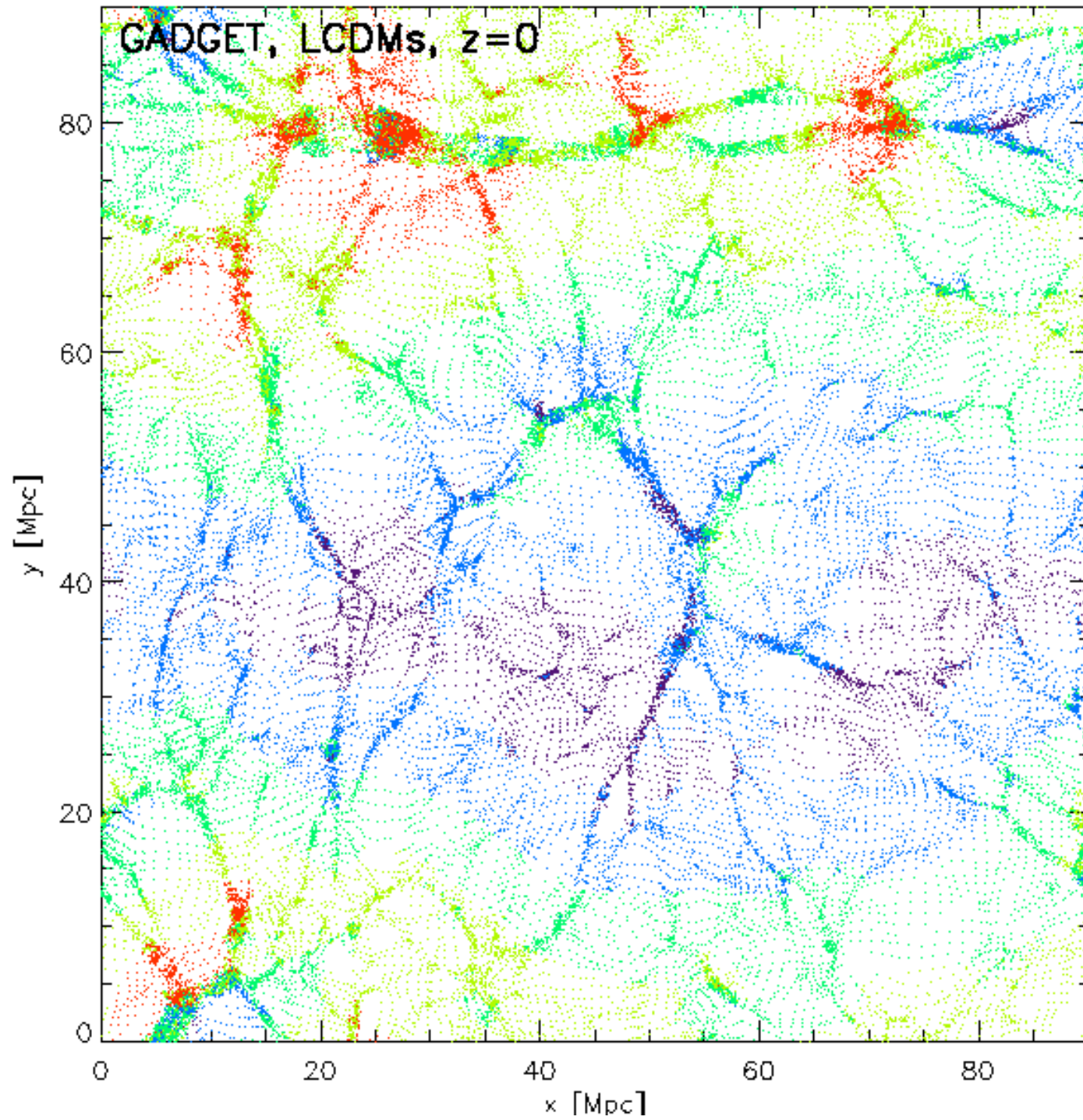
Answer: Not an accelerator physicist. Previous comments from a study on cosmology simulations



- **“Robustness of cosmological simulations I: Large-scale structure,” Heitmann, Ricker, Warren, Habib (astro-ph/0411795)**
- **Goal: test & compare 6 different N-body codes**
- **Uses 4 test problems**
- **Every code starts from exactly the same initial conditions**
- **Results are analyzed with the same set of analysis codes**
- **Investigation of particle 2-point functions, velocity statistics, halo catalogues**

Given the commonality of issues, close collaboration among computational cosmologists, plasma physicists, and accelerator physicists is natural & highly beneficial.



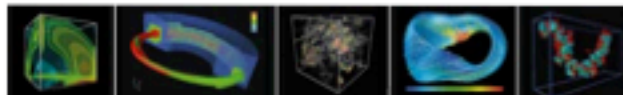


Signpost System 1985



Cray-2

- 244 MHz (4.1 nsec)
- 4 processors
- 1.95 Gflop/s peak
- 2 GB memory (256 MW)
- 1.2 Gflop/s LINPACK R_max
- 1.6 m² floor space
- 0.2 MW power



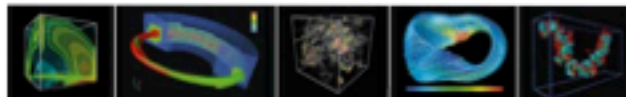
Signpost System in 2005

Artist's rendition of Blue Gene, a full-scale BG/L with 360 Tflop/s peak scheduled to become fully operational in early 2005. The computer's name is derived from its principle intended purpose — to model the folding of human proteins.



IBM BG/L @ LLNL

- 700 MHz (x 2.86)
- 65,536 nodes (x 16,384)
- 180 (360) Tflop/s peak (x 92,307)
- 32 TB memory (x 16,000)
- 135 Tflop/s LINPACK (x 110,000)
- 250 m² floor space (x 156)
- 1.8 MW power (x 9)



Preparing for the future



- High end platforms w/ 10's of thousands of procs are here now
 - Access to platforms w/ 100's of thousands of procs is coming
 - Petascale by the end of the decade
- The accelerator community needs to continue to foster close ties with the applied math/comp sci communities. Needs include:
 - Scalable solvers
 - Multiscale/multiresolution solvers
 - Optimization tools
 - Large-scale viz tools
 - Tools for managing massive amounts of data
 - High speed networks
 - Collaborative tools
 - And many others
- Using in concert with theory and expt, petascale computing will open new doors to understanding the physics of intense beams, designing EM structures, and developing new methods of particle acceleration.