Constraints on Structure Based Optical Accelerators

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High Brightness Beams Workshop Erice, Italy, October 2005

- Optical structures short overview
- Efficiency & Luminosity
- Non-linear phase-shift
- Heat dissipation
- Stress considerations

Constraints & Assumptions

- Gradient: order of 1[GV/m] or higher.
- *Efficiency*: limited by radiation source and acceleration scheme # Lasers anticipated efficiency of wall-plug to light 10% → 30%?!
 # Efficiency of acceleration scheme major *difficulty Breakdown*: at optical wavelengths *dielectrics* sustain higher
 - fields comparing to metals \rightarrow non-linear effects
- Manufacturing constraints favor planar structures consistent with luminosity constraints.
- Single Mode: width of vacuum tunnel $0.3\lambda 0.8\lambda \rightarrow position$
- *Machining tolerance*: 1µm at 3 cm wavelength. Four orders of magnitude difference are difficult to preserve at λ =1µm.

S. Banna, D. Schieber and L. Schächter; App. Phys. Lett., Vol.84(5) 723-5 (2004).
S. Banna, D. Schieber and L. Schächter; J. of Appl. Phys. 95(8) 4415-4426 (2004).

Structure parameters

$$Z_{\rm int} \equiv \left(E_0 \lambda\right)^2 / 2P$$

$$\begin{array}{c} \hline \textbf{Transverse PBG} \\ \hline \textbf{Transverse PBG} \\ \hline \textbf{Longitudinal PBG} \\ \hline \textbf$$

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L. Schächter; Energy Recovery in an Optical Linear Collider, Physical Review E **70**, 016504 (2004)



Single bunch & no feedback

•Wake parameter:

Decelerating field for a given charge

•Beam-loading parameter:

Beam-loading of the accelerating mode $E_{dec}^{(F)} \equiv \kappa_1 q$





$$\kappa_{1} = \frac{\gamma_{\rm gr}}{1 - \beta_{\rm gr}} \frac{\Delta_{\rm int}}{\sqrt{\mu_{0} / \varepsilon_{0}}} \frac{\pi}{4\pi\varepsilon_{0} / \varepsilon_{0}}$$

External Laser

 $E_{\rm dec} \equiv \kappa q$

Traveling-wave Acceleration Module

 $\kappa = 2 / 4\pi \varepsilon_0 R^2$

Contribution of the fundamental to the total deceleration

Stupakov & Bane, PR STAB, 2003



Single bunch & no feedback



Single bunch & no feedback

By splitting the bunch into a train of micro-bunches, the projection of the wake on the fundamental remains the same, but higher frequencies are suppressed



Train of bunches & no feedback

Laser Traveling-wave Acceleration Module

External

What is the efficiency in case of a train of micro-bunches?

For an answer, one needs to make two observations:

a) The laser pulse duration ought to be longer in order to account for the <u>macro-bunch length</u>.

b) The envelope of the laser pulse must be **tapered**, in order to compensate for <u>beam loading</u>.



Train of bunches & no feedback



In spite of splitting the macrobunch, there still is 50% waste of the laser energy \rightarrow feedback loop.

External

Traveling-wave Acceleration Module

0.0

Laser









Luminosity

Fluence

	NLC	Single no FB	Single no FB	Train & FB	Train & FB
<i>f</i> _{rr} [Hz]	120	1(6)	1(9)	1(6)	1(6)
N _{micro}	190	1	1	1(3)	1(3)
N _{el}	0.75(10)	5(4)	5(4)	1(6)	1(6)
σ _z [nm]	0.11(6)	4	4	4	4 Density !
σ _x [nm]	245	245	24.5	245	24.5
σ _y [nm]	2.7	2.7	0.27	2.7	0.27
n ^{-1/3} [Å]	2.1	3.8	0.8	1.4	0.3
H _D	1.4	1.0	1.0	1.0	1.0
L[cm ⁻² sec ⁻¹]	2(34)	3(25)	3(30)	1.2(34)	1.2(36)
Beam power [MW]	13.7	4(-3)	4	80	80

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Polarization non-linearity

- Breakdown & Nonlinearity: In dielectrics at optical wavelengths breakdown threshold is higher than that required to excite nonlinear effects.
- ◆ Main Nonlinear Effect: laser affects dielectric coefficient
 → phase shift.

 $\vec{P} = \varepsilon_0 \left[\chi^{(1)} \vec{E} + \chi^{(2)} \vec{E} \vec{E} + \chi^{(3)} \vec{E} \vec{E} \vec{E} + \dots \right]$

- Third harmonic generation
- Four wave mixing
- Nonlinear refraction

- Second harmonic generation
- Sum-frequency generation

$$\begin{array}{l} \textbf{Nonlinear Phase-shift} \\ \textbf{Nonlinear Phase-shift} \\ \textbf{Planar Structure} \\ \textbf{D}_{max} \\$$

Non-linear Phase-shift criterion

 $\sigma_z \leftrightarrow$ $\sigma_z = \frac{\alpha}{2\pi} \lambda$ $0 < \alpha < \frac{2\pi}{12}$ Phase-shift less than 10% Efficient acceleration Synchronization $\Delta\phi_{NL} < \frac{2\pi}{120}$ Condition

Solution: Redesigning the structure to (initially) have a phase velocity larger than *c*.

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Heat flow – single pulse operation

$$\Delta T(r = D_{\text{ext}}) = \frac{1}{h} \left(\frac{\partial}{\partial x} \Delta T \right]_{x = D_{\text{ext}}}$$

$$\left[\frac{\partial}{\partial t} - \frac{1}{D}\nabla^2\right]\Delta T = \frac{1}{D\sigma_T} p_{\text{loss}}(r, z, t)$$

Diffusion Coefficient $D[\sec/m^2]$ Heat Conductivity $\sigma_T[J/\sec m^{\circ}K]$ $\left(\frac{\partial}{\partial x}\Delta T\right]_{x=D_{\text{int}}} = 0$

Dissipated EM Power Density

$$p_{\text{loss}}(r,z,t) \simeq \varepsilon \tan \delta \frac{\omega_0}{c} \langle S_z(r,z,t) \rangle$$

Heat Flow - Train of Pulses
$$T_{rr} > 10[nsec]$$

 $T_p \sim 1[psec]$ $\Delta T(x,z;t) = \frac{P_0}{\sigma_T} \frac{T_p}{T_{rr}} \sum_{v=-\infty}^{\infty} sinc \left(\pi v \frac{T_p}{T_{rr}}\right) exp \left(j2\pi v \frac{t-z/V_{gr}}{T_{rr}}\right)$
Power profile $\Delta T(x,z;t) = \frac{P_0}{\sigma_T} \frac{T_p}{T_{rr}} \sum_{v=-\infty}^{\infty} sinc \left(\pi v \frac{T_p}{T_{rr}}\right) exp \left(j2\pi v \frac{t-z/V_{gr}}{T_{rr}}\right)$
Power profileFor the parameters of interest,
the deviation from average
 ΔT is negligible !! $D_{ext} - D_{int}$ $D_{$



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EM Stress on a Planar Accelerator

Total repelling Laser Pressure – Time-averaged Maxwell stress-tensor component

Time-averaged Lorentz, force density inside the v'th layer is due to the polariz,ation current

Total time-averaged boundary stress between the v'th and the v+1'th layer is due to polarization surface charge

$$\langle T_{xx} \rangle = \frac{1}{4} \varepsilon_0 |E_x|^2 - \frac{1}{4} \varepsilon_0 |E_z|^2 - \frac{1}{4} \mu_0 |H_y|^2$$
$$= -\frac{1}{4} \varepsilon_0 |E_0|^2$$
$$E_0 = 1 \left[\frac{\text{GV}}{\text{m}} \right] \Rightarrow |\langle T_{xx} \rangle| \approx 2.1 \left[\frac{\mu \text{N}}{\mu \text{m}^2} \right]$$
$$\approx 2 \text{[MP]}$$

≪1[GP]

$$f_x = \frac{1}{2} \operatorname{Re} \left\{ j \omega \varepsilon_0 \left(\varepsilon_v - 1 \right) E_z^* \mu_0 H_y \right\}$$

$$F_{x} = \frac{1}{4}\varepsilon_{0} \left| E_{x,\nu}^{2} \right| \left[\frac{\varepsilon_{\nu}^{2}}{\varepsilon_{\nu+1}^{2}} - 1 \right]$$





Summary: Dielectric Structures Longitudinal PBG **Transverse PBG** Dielectric periodic structures may confine an accelerating mode allowing high order modes to leak out. Cylindrical Bragg Structure Planar Bragg Structure

Train of micro-bunches contributes to suppression of high frequency wakes.

Summary: Configurations



Summary: Efficiency & Luminosity



 (I) Amplifier compensates for all radiation loss (active enhancement of the Q—factor) facilitated by the wake being "quasi-coherent"
 (II) External laser compensates for beam-loading (tapered pulse)
 (III) Luminosity

Summary

- (I) With commonly used materials Kerr effect may become a problem. By accounting for the local increase in the dielectric coefficient, it is possible to compensate for the phase velocity reduction.
- (II) Preliminary estimates indicate that heat dissipation needs close attention because the high repetition rate dictated by the luminosity constraint. Employing optical fiber technology may eliminate the problem although the latter works well in homogeneous medium, implications on multilayered structure are unclear as yet.

(III) The typical stress is evaluated to be of the order of 1[MPascal] which is three orders of magnitude below typical critical values at μm scale.