# Harnessing plasma waves as radiation and particle sources Dino Jaroszynski

#### **University of Strathclyde**

## Outline of talk

- How can we use a laser-driven plasma wave to produce coherent electromagnetic radiation?
- Laser driven wakes
- Alpha-X
- Ultra-short bunch electron production using wakefield accelerators
- Synchrotron and free-electron laser sources

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Conclusion and outlook





### Wakefield acceleration





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# Radiation sources: Synchrotron and Free-electron laser (FEL)

- Use output of wakefield accelerator to drive compact synchrotron light source or FEL
- Take advantage of electron beam properties
- Coherent spontaneous emission: prebunched FEL  $I \sim I_0(N+N(N-1)f(k))$
- Operate in superradiant regime: FEL x-ray amplifier (self-similar evolution)

Potential compact future synchrotron source and x-ray FEL

- Need a low emittance GeV beam with < 50 fs electron beam with I > 1 kA
- Operate in superradiant regime: SASE alone is not adequate
- Need to consider injection or pre-bunching





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# ALPHA-X Programme

#### Main areas of research:



Advanced Laser Plasma High-energy Accelerators towards X-rays: ALPHA-X (a-x)

#### Collaborative project involving groups from the UK, EU and US











- Strathclyde injector, laser-plasma & FEL: experiments & theory
- CCLRC RAL theory & exps.: wakefield studies and diagnostics
- Oxford plasma channels
- Imperial all-optical injector, laser-plasma acceleration
- CCLRC Daresbury Injector, undulator & FEL
- Dundee injector, electron diagnostics & FEL
- St Andrews University theory
- GOALS:Accelerate to 1 GeV in 1cm using a wakefield accelerator.<br/>Demonstrate laser-driven light source: synchrotron or FEL

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### People involved in project

• Dino Jaroszynski, Klaas Wynne, Bob Bingham, Ken Ledingham, Albert Reitsma, Yuri Saviliev, Slava Pavlov, Riju Issac, David Jones, Bernhard Ersfeld, Steven Jamison, Jordan Gallcher, Andrey Lyachev, Enrico Brunetti, Mark Wiggins – *Strathclyde* 

• Karl Krushelnick , Bucker Dangor , Zulfika Najmudin, Malte Kaluza, Alex Thomson, Stuart Mangles – Imperial College

• Bob Bingham, Henry Hutchinson, Peter Norreys, Raoul Trines, Kate Lancaster, Chris Murphy – RAL CCLRC

• Simon Hooker, Keith Burnett, Ian Walmsley, Justin Wark, Tony Gonsalves – Oxford

• Allan Gillespie, Allan McCloud, Steven Jamison, – Dundee, Abertay-Dundee

- Alan Cairns St Andrews
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Strathclyde Electron and Terahertz to Optical Pulse Source Gennady Shvets – Austin Texas

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- Fred van Goor, Arsen Khachatryan, Kees van der Geer, Marieke Loos, Bas van der Geer, van der Wiel – *The Netherlands*
- Andrey Savilov, Vladimir Bratman IAP, Nizhniy Novgorod

Total of 63 people

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### Sources: TOPS ASTRA, and OXFORD

TOPS (Strathclyde): 5 TW source (800nm, 50fs 10Hz 250mJ) upgrade to 1.1J (20 TW) October 2005 RAL: ASTRA 1 J, 20 TW source and Gemini: 0.5 PW (in 2007) Oxford: 2 TW source Strathclyde: 8 MeV High-brightness sub-picosecond photoinjector – being constructed Other European facilities: FELIX, Lund, MPQ, DESY





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Compact femtosecond duration synchrotron or FEL source

 $\lambda$  = 2.8 nm – 1 mm (<1GeV beam)







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# Capillary: preformed plasma waveguide

	$r^2$	$r_0 = 150 \ \mu m \ capillary$
	$n(r) = n_0 + \delta n \frac{1}{r_0^2}$	<i>n</i> (0) = 10 <sup>18</sup> cm <sup>-3</sup>
	$\frac{1}{1}$	/ > 10 <sup>17</sup> W/cm <sup>2</sup>
	$w_m = \left(\frac{r_0^2}{\pi r_e \delta n}\right)^4$	Plasma formed by electrical discharge
electrodes		between electrodes
	Plasma capillary waveguide	
phase-m	2 – 5 cm atching: tapered ca	pillary
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### Plasma waveguide formation

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- After t ~ 80 ns plasma in quasi equilibrium.
- Ohmic heating of plasma balanced by conduction of heat to wall

$$\frac{1}{r}\frac{d}{dr}\left(r\kappa_{e\perp}\frac{dT}{dr}\right) + \sigma_{\perp}E^{2} = 0$$

Solution of the heat flow equation yields scaling relation for matched spot size:

$$N_{M}$$
[µm] = 1.5 × 10<sup>5</sup>  $\frac{\sqrt{a}$ [µm]}{( $\overline{N}_{e}$ [cm<sup>-3</sup>])<sup>1/4</sup>

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(See S. Hooker et al)

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# **Injection: Promising results**



- Many groups have now observed monoenergetic beams
- Still need to demonstrate controlled acceleration
- Measurement of electron bunch duration needed to determine peak current



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### ALPHA-X all-optical injection experiments on ASTRA





### Beam stability: electron beam pointing



#### Krushelnick et al. Imperial College

- Electron beam profile for E > 11 MeV
- Pointing instability ~ 3°
- Multiple beamlets observed.

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- Lanex screen 480mm from target
- Contours at 1°, 2°, 3°, 4°

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ource

### Beams have very narrow divergence





### LOA – Malka





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### LBNL Leemans

### 2004

- Channel guided wakefield accelerator
- 2.4 mm gas jet
- *n* = 4.5 x 10<sup>19</sup> cm<sup>-3</sup>
- *E* = 500 mJ; 55 fs FWHM; *w* = 8.5 μm
- f/4 off axis parabola
- / = 1.1 x 10<sup>19</sup> Wcm<sup>-2</sup>

C. Geddes et al. Nature 2004





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### State-of-the-art laser driven wakefield accelerators

- Electron bunch charge of the order of 22 pC in the RAL experiments.
- In contrast, the LOA (gas jet) and LBNL (channel) experiments produced 500 pC, 170 MeV and 160 pC, 100 MeV bunches, at n<sub>e</sub> = 7 × 10<sup>18</sup> cm<sup>-3</sup> and n<sub>e</sub> = 2 × 10<sup>19</sup> cm<sup>-3</sup>, respectively.
- Estimated e-bunch length in these experiments few femtoseconds
- Peak current should be multi-kilo-Amperes (~ 20 kA for 5 fs 100 pC).
- Emittance was estimated to be of the order of  $\varepsilon \sim 1 \pi$  mm mrad.
- The energy spread ~ 1 3 % suitable for FEL applications with more work on emittance and energy spread (need to get to 0.5% for FEL and linear collider stage)
- 6.5 MeV 100 fs photoinjector (BNL/Eindhoven/LAL gun)





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#### Modelling of Laser Wakefield Acceleration





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### Laser pulse envelope dynamics Strathclyde

laser pulse amplitude: a<sub>0</sub>

laser pulse energy depletion rate:  $\omega_d \sim a_0^2 \omega_s$ 



**Linear regime:**  $a_0^2 \ll 1$ ,  $\omega_d \ll \omega_s$ : pulse energy loss through photon deceleration without envelope modulation, static wakefield, low energy efficiency

**Nonlinear regime:**  $a_0^2 \sim 1$ ,  $\omega_d \sim \omega_s$ : pulse energy loss through photon deceleration and strong envelope modulation, dynamic wakefield, better energy efficiency

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Contracting of Contra

## Beam line



### Transport of flat and curved femtosecond bunches: buncher

Nps=200,000, Rb=1.8 mm, sc3D, rfsol=0.478, plasmasol=0.927 -> 46 FWHM  $\mu m,$  331 FWHM fs, 1.73  $\mu m,$ 





Figure 3-7: GPT simulations of a 10 pC bunch at entrance of the plasma accelerator, initial radius is 5 mm

Nps=200,000, Rb=5 mm, sc3D -> 34 FWHM µm, 107 FWHM fs, 3.1 µm.

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# Electron distribution at capillary entrance



#### van der Geer et al. 2005



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### Electron bunch evolution: bunch compressor

Implications for wakefield accelerators:

electron bunches cannot be expanded radially:

beam will expand by 6 fs if expanded to 2 mm in 0.5 m

 $\delta t = D^2/4fc$ 



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van der Geer et al. 2005



### Injection-compression schemes

to minimize energy spread, bunch duration *short* compared to *plasma period*very sensitive to *timing jitter* between laser pulse and electron bunch
an interesting option: injection *in front of* laser pulse

initial electron energy low →
 bunch *slips backwards* through laser pulse

- electrons trapped behind laser pulse
- wakefield acts as *relativistic mirror* and accelerates electrons to *high energy*

• scheme depends on *overlap* of focusing and accelerating regions for *nonlinear* or *narrow channel* wakefields







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# Trapping and acceleration



#### example:

z-vt

(*z*-*vt*)

- 4 MeV beam in 3.5 x 10<sup>17</sup> cm<sup>-3</sup> plasma
- Laser pulse acts as filter, rejecting electrons outside spot size
- Initial bunch spot size can be *larger* than laser spot size (i.e. if we accept losses)
- Contrast with conventional injection requires small bunch spot to avoid energy spread
- Accelerated bunch has low transverse emittance

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# Linear plasma model



GPT: Loose and van der Geer (Pulsar)
Electron bunch longer than plasma period
Quasistatic plasma model (Albert Reitsma)



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### Undulator: compact synchrotron source and free-electron laser

- Slotted pole focussing undulator (Daresbury) Strathclyde/Daresbury
- 200 period
- Undulator period 1.5 cm
- Minimum gap 5 mm
- Undulator parameter a<sub>u</sub> ~ 1



undulator magnet

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# Compact synchrotron source





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## ALPHA-X synchrotron source

- Radiation wavelength: 0.2 nm (1 GeV beam)
- Pulse length depends on Δγ/γ ( ~ 2% measured) and electron bunch duration
- Ideal  $\Delta\lambda/\lambda = 1/2N_u = 0.5\%$
- Ideal pulse length: sub-femtosecond (point charge)
- Synchrotron radiation pulse length ~ 10 fs



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# Energy spread and emittance

- Non-ideal electron beams degrade radiation properties: reduces brilliance and lowers FEL gain
- Emittance: area in transverse phase space

$$\varepsilon_{x,y} \approx \sigma_{x,y} \sigma'_{x,y}$$

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- Approx. beam size times RMS beam divergence
- Normalised emittance:  $\varepsilon_n = \gamma \beta \varepsilon$  needs to be  $\approx 1 \pi mm$ mrad
- Energy spread:

needs to be < 
$$\rho$$
 < 0.5%

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### Predicted synchrotron radiation peak brilliance

 $I(k) \sim I_0(k)(N+N(N-1)f(k))$ 

- $N_u = 200$ 
  - $= 1 \pi \text{ mm mrad}$ 
    - = 10 fs

Q = 100 – 200 pC FEL: Brilliance 5 – 7 orders of magnitude larger





ε<sub>n</sub>

 $\tau_{e}$ 

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### Free-electron laser driven by wakefield accelerator

- Combination of undulator and radiation fields produces ponderomotive force which bunches electron on a wavelength scale
- Laser field grows exponentially at a rate governed by
- Matched electron beam:  $\rho = 1.1 \gamma^{-1} B_u \lambda_u^{4/3} I_{pk}^{1/3} \varepsilon_n^{-1/3}$
- Gain length
- Need  $\delta \gamma / \gamma < \rho$  and  $\varepsilon_n < 4\lambda \beta \gamma \rho / \lambda_u$  or  $\varepsilon_n < \gamma \lambda$  (matched) See Bonifacio et al. (Nuova Cimento 1990, 1992)

 $L_g = \frac{\lambda_u}{2\pi\sqrt{3\rho}}$ 



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## Free-electron laser

**Ultimate goal of technology:**  $\lambda \sim 1.5$  nm x-ray FEL requiring  $\sim 1$  GeV

- Growth of an injected or spontaneous field in a FEL amplifier is given by  $I = I_0 \exp(gz)$ ,
- $g = 4\pi\rho 3^{1/2}/\lambda_u \text{small signal gain}$ ,
- $\rho$  FEL gain parameter a function of the beam energy, current and emittance.
- $\varepsilon_n = \gamma \sigma \Omega$  is the normalised emittance of the beam.
- $\rho \sim 0.001$  to 0.02 for our expected electron beam parameters but need energy spread  $\delta\gamma / \gamma < 2\rho$  i.e. 0.2 4%
- •gain length < 10  $\lambda_u$

#### Makes sense to optimise gain, if possible

*EXAMPLE:* ALPHA-X: 4 nm source assuming a 1 GeV beam with 100 pC charge and a duration of 10 fs we get a peak current of 10 kA. With a 1.5 cm period undulator with a field of B field of 1 T we get a  $\rho \approx 0.005$ , which gives a gain length of  $10\lambda_u$  and a constraint on the energy spread of  $\approx 1\%$ , which may be achievable. To achieve saturation we need about 100 - 200 undulator periods.



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# Conclusion

- Femtosecond compact radiation sources feasible
- Compact synchrotron
- Compact FEL
- Large tuning range THz to x-ray (water window)
- Challenges: emittance, peak current, energy spread
- More challenges: synchronism, shot to shot jitter, bunch duration
- Outlook good with a lot of work need a roadmap
- Linear collider application far away. Need to develop laser technology and schemes for staging





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