

E & M memorization sheet

①

Poisson's eq. $\nabla \cdot \vec{E} = 4\pi \rho$

Faraday's law: $\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$

moving charge: $\vec{v} = \frac{c^2 \times \vec{B}}{B^2}$, $\vec{v}_{drift} = c \frac{\vec{E} \times \vec{B}}{B^2}$

no name: $\nabla \cdot \vec{B} = 0$

Ampère's law: $\nabla \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$

} spatial and temporal variables are not set by Maxwell's eqs

structure of Conservation Laws:

$$\frac{\partial \Psi}{\partial t} + \nabla \cdot \vec{J} = S$$

Ψ - scalar quantity of interest... 'stuff'
 \vec{J} - current of stuff (flow of Ψ)
 S - local source or sink of stuff

conservation of charge: $\frac{\partial \rho}{\partial t} + \nabla \cdot \vec{j} = 0$

conservation of Energy: $\frac{\partial}{\partial t} \left[\underbrace{\frac{1}{8\pi} (\vec{E}^2 + \vec{B}^2)}_{\text{energy}} \right] + \nabla \cdot \left[\underbrace{\frac{c}{4\pi} (\vec{E} \times \vec{B})}_{\text{Poynting vector}} \right] = \underbrace{-\vec{E} \cdot \vec{j}}_{\text{ohmic power}} = -\frac{d}{dt} \left(\frac{1}{2} \rho_m v^2 \right) \equiv U_m$

$$\Rightarrow \frac{d}{dt} [U_E + U_B + U_m] + \nabla \cdot \vec{S} = 0$$

$$(\nabla \times \vec{B}) \times \vec{B} = \underbrace{-\nabla \left(\frac{|\vec{B}|^2}{2} \right)}_{\text{pressure (compression)}} + \underbrace{(\vec{B} \cdot \nabla) \vec{B}}_{\text{tension (compression)}}$$

conservation of momentum: $\frac{d}{dt} (\rho_m \vec{v}) + \frac{\partial}{\partial t} \left[\underbrace{\frac{\vec{S}}{c^2}}_{\text{e/m momentum density}} \right] = \nabla \cdot \underbrace{\vec{T}}_{\text{force per unit volume}}$

$$\Rightarrow \frac{d}{dt} \left[\underbrace{\vec{p}_{mech}}_{\rho_m \vec{v}} + \underbrace{\vec{p}_{em}}_{\frac{1}{4\pi c} (\vec{E} \times \vec{B})} \right] = \underbrace{\vec{f}_{em}}_{\nabla \cdot \vec{f}}$$

$$\cdot \vec{T}_{op} = \frac{1}{4\pi} [E_x E_x + B_x B_x - \frac{1}{2} \delta_{\alpha\beta} (E^2 + B^2)]$$

vector wave equation: $\frac{\partial^2 \vec{E}}{\partial t^2} - c^2 \nabla^2 \vec{E} = 4\pi \frac{\partial \vec{j}}{\partial t} + 4\pi c^2 (\nabla \rho)$

$$\frac{\partial^2 \vec{B}}{\partial t^2} - c^2 \nabla^2 \vec{B} = 4\pi c (\nabla \times \vec{j})$$

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$\frac{\partial^2 \vec{A}}{\partial t^2} - c^2 \nabla^2 \vec{A} = 4\pi c \vec{j} - c \nabla \left[\frac{\partial \phi}{\partial t} + c \nabla \cdot \vec{A} \right]$$

$$\nabla^2 \phi = -4\pi \rho - \frac{1}{c} \frac{\partial}{\partial t} \nabla \cdot \vec{A}$$

Coulomb gauge: $\nabla \cdot \vec{A} = 0$

Lorentz gauge: $\frac{\partial \phi}{\partial t} + c \nabla \cdot \vec{A} = 0$

canonical momentum eq:

$$\frac{d}{dt} \left[\underbrace{\vec{p}}_{\equiv \vec{p}_{\text{canonical}}} + \frac{q}{c} \vec{A} \right] = -q \nabla \left[\underbrace{\phi - \frac{\vec{v} \cdot \vec{A}}{c}}_{\equiv \phi_{\text{effective}}} \right] \Rightarrow H = \frac{(\vec{p}_{\text{can}} - \frac{q}{c} \vec{A})^2}{2m} + q \phi$$

3 basic properties of electric dipole moment: (1) energy, (2) force, and (3) torque:

1) $U = -\vec{p} \cdot \vec{E}$

2) $\vec{F} = \vec{p} \cdot \nabla \vec{E}$

3) $\vec{\tau} = \vec{p} \times \vec{E}$

For magnetic systems

$$U = -\vec{m} \cdot \vec{B}$$

$$\vec{F} = \vec{\nabla} \cdot \vec{B}$$

$$\vec{\tau} = \vec{m} \times \vec{B}$$

Boundary conditions: $\Delta E_{||} = 0$ $\Delta E_{\perp} = 4\pi \sigma_{\text{total}}$

$\Delta D_{\perp} = 4\pi \sigma_{\text{free}}$ $\Delta P_{\perp} = -\sigma_{\text{bound or polarization}}$

$$\vec{H} = \vec{B} - 4\pi \vec{M}, \quad \nabla \times \vec{H} = \frac{4\pi}{c} \vec{j}_{\text{conduction}}, \quad \vec{B} = \vec{\mu} \cdot \vec{H}$$

Bessel's eqn: $\frac{1}{r} \frac{d}{dr} \left(r \frac{d}{dr} f \right) + \left(k^2 - \frac{\nu^2}{r^2} \right) f = 0$

Soln. $J_{\nu}, Y_{\nu} \rightarrow$ standing waves

$H_{\nu}^{(1)}, H_{\nu}^{(2)} \rightarrow$ travelling waves

Green's Function in cylindrical coordinates

$$\left. \begin{matrix} G_>(p;p') \\ G_<(p;p') \end{matrix} \right\} = \sum_{m=-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{dk}{(2\pi)^2} e^{im(\theta-\theta')} e^{ik(z-z')} \quad (4\pi) \quad \left\{ \begin{matrix} K_m(k\rho) I_m(k\rho') \\ I_m(k\rho) K_m(k\rho') \end{matrix} \right.$$

note: spherical solution to Laplace's eqn: Legendre Poly has the orthogonal relation

$$\int_{-1}^{+1} dx P_l(x) P_{l'}(x) = \frac{2}{2l+1} \delta_{l,l'}$$

if problem is symmetric only even $P_l(x)$ are non-zero, if antisymmetric only odd P_l are non zero.

$$P_l(x) = \begin{cases} \text{even } l & \text{even} \\ \text{odd } l & \text{odd} \end{cases}$$

$$\Phi(r, \theta) = \sum_{l=0}^{\infty} \left[a_l r^l + \frac{b_l}{r^{l+1}} \right] P_l(\cos \theta)$$

→ single charges are diamagnetic

→ magnetization current density: $\vec{j}_m = c \nabla \times \vec{M}$

bound charge: $\rho_b = -\nabla \cdot \vec{P}$

+ surface current: $\vec{j}_s = c \vec{M} \times \vec{n}$

surface charge: $\sigma_b = \vec{P} \cdot \vec{n}$

$$u_{\text{total}} = \frac{\vec{H} \cdot \vec{B}}{8\pi}, \quad \frac{\vec{D} \cdot \vec{E}}{8\pi}, \quad \vec{D} = \vec{E} + 4\pi \vec{P}$$

Force density $\frac{d\vec{F}}{dV} = \vec{M} \cdot \nabla \vec{B} = \left(\frac{\mu}{4\pi} \right) \vec{H} \cdot \nabla \vec{B}, \quad \vec{M} = \frac{\chi - 1}{4\pi} \vec{H}$

→ $K = \pm \frac{\omega}{c}$ in vacuum, $K = \pm \frac{\omega}{c} \sqrt{\epsilon(\omega)}$ otherwise

$$\vec{S} = \frac{c}{4\pi} \vec{\nabla} \times \vec{B}^*$$

→ circulating Poynting flux (attached to ob, etc.)

$$\vec{P} = \int \vec{E} \cdot \vec{E}^*$$

→ 'power' mechanical energy

$$\nabla \cdot (\text{Im } \vec{S}) = -\text{Im } \vec{P} + \frac{A_0}{4\pi} [|\vec{B}|^2 - |\vec{E}|^2]$$

dielectric coefficient

$$\epsilon(\omega) = 1 + \frac{4\pi N \alpha(\omega)}{1 - \rho N \alpha(\omega)}, \quad \alpha(\omega) = \sum_l \frac{F_l}{\omega_l^2 - \omega^2 - i\gamma_l \omega}$$

$$V_g(\omega) = \frac{V_p(\omega)}{1 + \omega \frac{R}{\omega} \ln \frac{1}{\epsilon_R}}$$

From Faraday's law can get relationship between \vec{E} & \vec{B} $\therefore \vec{K} \times \vec{E} = \frac{1}{c} \frac{d\vec{B}}{dt}$

$$\frac{|\vec{B}|^2}{|\vec{E}|^2} = \epsilon_R \left[1 + \frac{1}{4} \left(\frac{\epsilon_I}{\epsilon_R} \right)^2 \right]$$

AC conductance $\rightarrow \omega \epsilon_I$

AC reactance $\rightarrow \omega(1 - \epsilon_R)$

For conducting media: $\epsilon = 1 + i \frac{4\pi\sigma}{\omega}$

time averaged power: $\langle dP \rangle = r^2 d\Omega \langle \vec{S} \cdot \hat{r} \rangle$

per solid angle $\langle \frac{dP}{d\Omega} \rangle = r^2 \langle \vec{S} \cdot \hat{r} \rangle$

radiation resistance of antenna: $R_{rad} = \frac{\int d\Omega \langle \frac{dP}{d\Omega} \rangle}{\frac{1}{2} I^2}$

Liénard-Wiechert Potentials:

$\vec{A}(\vec{r}, t) = q \left[\frac{\vec{\beta}}{R - \vec{R} \cdot \vec{\beta}} \right]_{retarded\ time}$

$\phi(\vec{r}, t) = q \left[\frac{1}{R - \vec{R} \cdot \vec{\beta}} \right]_{retarded\ time}$

$\vec{S}_a = \frac{q^2}{4\pi c^3 R^2} |\dot{\vec{a}}|^2 \sin^2 \theta \hat{R}$

Larmor radiation formula: $P = \frac{2}{3} \frac{q^2}{c^3} |\dot{\vec{a}}|^2$

$\dot{\vec{a}} = \frac{\ddot{\vec{p}}}{q}$

$\langle P \rangle = \frac{\omega^4}{3c^3} |\dot{\vec{p}}|^2$

