

General useful stuff

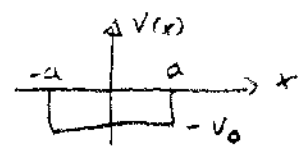
$$\sum_{n=0}^{\infty} x^n = \frac{1}{1-x} \quad (1)$$

binomial series $(1+x)^n = 1 + nx + \frac{n(n-1)}{2!} x^2 + \dots$

Taylor series (about a) $f(a+x) = f(a) + x f'(a) + \frac{x^2}{2!} f''(a) + \dots$

limits $(1 + \frac{x}{n})^n \rightarrow e^x$ as $n \rightarrow \infty$; $\frac{x^n}{n!} \rightarrow 0$ as $n \rightarrow \infty$
 $x \ln x \rightarrow 0$ as $x \rightarrow 0$; $n^c x^n \rightarrow 0$ as $n \rightarrow \infty$ if $|x| < 1$

quantum perturbed eigenvalues: $E'_k = E_k + \langle \psi_k | \hat{H}' | \psi_k \rangle + \sum_{n \neq k} \frac{|\langle \psi_k | \hat{H}' | \psi_n \rangle|^2}{E_k - E_n} + \dots$

Potential well:  bound states $V_0 < E < 0$
 $\tan qa = \begin{cases} k/|q| & \text{even parity} \\ -q/|k| & \text{odd parity} \end{cases}$ $q^2 - |k|^2 = \frac{2m|V_0|}{\hbar^2}$

particles in a rectangular box: Energy levels: $E_{lmn} = \frac{\hbar^2}{8m} \left(\frac{l^2}{a^2} + \frac{m^2}{b^2} + \frac{n^2}{c^2} \right)$ $l, m, n \geq 1$
 $M = \text{mass}$

Thermo:

1st law $dU = dQ + dW$
 ↑ ↑ ↑
 change in internal energy work done on system heat supplied to system

2nd law $dS \geq \frac{dQ}{T}$

Thermodynamic work: $dW = -p dV = \gamma dA = \vec{E} \cdot d\vec{p} = \vec{B} \cdot d\vec{m} = d\phi dq$
 ↑ ↑ ↑ ↑ ↑
 surface tension induced electric dipole moment magnetic dipole moment potential difference charge moved

Cycle coefficients:

Heat engine: $\eta = \frac{\text{work extracted}}{\text{heat input}} \leq \frac{T_h - T_c}{T_h}$

Refrigerator: $\eta = \frac{\text{heat extracted}}{\text{work done}} \leq \frac{T_c}{T_h - T_c}$

heat pump: $\eta = \frac{\text{heat supplied}}{\text{work done}} \leq \frac{T_h}{T_h - T_c}$

Heat capacities:

$C_V = \left(\frac{dQ}{dT} \right)_V = \left(\frac{\partial U}{\partial T} \right)_V = T \left(\frac{\partial S}{\partial T} \right)_V$; $\gamma = \frac{C_P}{C_V}$

$C_P = \left(\frac{dQ}{dT} \right)_P = T \left(\frac{\partial S}{\partial T} \right)_P$; $C_P - C_V = \left[\left(\frac{\partial U}{\partial V} \right)_T + P \right] \left(\frac{\partial V}{\partial T} \right)_P$

Thermodynamic potentials:

$$dU = TdS - pdV + \mu dN$$

enthalpy: $H = U + pV$ $dH = TdS + Vdp + \mu dN$

Helmholtz free energy: $F = U - TS$ $dF = -SdT - pdV + \mu dN$

Gibbs free energy: $G = U - TS + pV$ $dG = -SdT + Vdp + \mu dN$

Maxwell's relations

$$\left(\frac{\partial T}{\partial V}\right)_S = -\left(\frac{\partial p}{\partial S}\right)_V = \frac{\partial^2 U}{\partial S \partial V} ; \left(\frac{\partial T}{\partial p}\right)_S = \left(\frac{\partial V}{\partial S}\right)_p = \frac{\partial^2 H}{\partial p \partial S}$$

$$\left(\frac{\partial p}{\partial T}\right)_V = \left(\frac{\partial S}{\partial V}\right)_T = \frac{\partial^2 F}{\partial T \partial V} ; \left(\frac{\partial V}{\partial T}\right)_p = -\left(\frac{\partial S}{\partial p}\right)_T = \frac{\partial^2 G}{\partial T \partial p}$$

Latent heat: $L = T(S_2 - S_1)$

Clausius-Clapeyron eq. $\frac{dp}{dT} = \frac{S_2 - S_1}{V_2 - V_1} = \frac{L}{T(V_2 - V_1)}$

ideal gas

Adiabatic equations: $pV^\gamma = \text{const}$, $TV^{(\gamma-1)} = \text{const}$, $T^\gamma p^{(1-\gamma)} = \text{const}$

internal energy: $U = \frac{nRT}{\gamma-1}$

Virial expansion: $pV = RT \left(1 + \frac{B_2(T)}{V} + \frac{B_3(T)}{V^2} \right)$

\downarrow virial coef \downarrow virial coef
 $B_2(T)$ $B_3(T)$

Monatomic gas: $U = \frac{3}{2} NkT$, p pressure = $\frac{1}{3} nm \langle c^2 \rangle$ ← mean squared particle velocity

$$C_V = \frac{3}{2} Nk, C_P = C_V + Nk = \frac{5}{2} Nk, \gamma = \frac{5}{3}$$

$$S = Nk \ln \left[\left(\frac{m k T}{2 \pi \hbar^2} \right)^{3/2} e^{5/2} \frac{V}{N} \right]$$

Maxwell-Boltzmann distribution

particle energy distribution: $p(E)dE = \frac{2E^{1/2}}{\pi^{1/2} (kT)^{3/2}} e^{-\beta E} dE$

mean speed: $\langle c \rangle = \left(\frac{8kT}{\pi m} \right)^{1/2}$

Class partition

Classical: $\bar{E}_q = \frac{1}{2} kT$

ideal gas heat capacities: $C_v = \frac{1}{2} f N k$, $C_p = Nk \left(1 + \frac{f}{2} \right)$, $\gamma = \frac{C_p}{C_v} = 1 + \frac{2}{f}$
of degrees of freedom

Macroscopic thermodynamic variables

$$F = -kT \ln Z$$

$$U = F + TS = - \left(\frac{\partial \ln Z}{\partial \beta} \right)_{V, N}$$

$$S = - \left(\frac{\partial F}{\partial T} \right)_{V, N} = \frac{\partial (kT \ln Z)}{\partial T} \Big|_{V, N}$$

$$P = - \left(\frac{\partial F}{\partial V} \right)_{T, N} = - \left(\frac{\partial (kT \ln Z)}{\partial V} \right) \Big|_{T, N}$$

$$\mu = \left(\frac{\partial F}{\partial N} \right)_{V, T} = - \frac{\partial (kT \ln Z)}{\partial N} \Big|_{V, T}$$

Occupation #

$\langle n_i \rangle = \frac{1}{e^{\beta(\epsilon_i - \mu)} - 1}$ Bose-Einstein

$\langle n_i \rangle = \frac{1}{e^{\beta(\epsilon_i - \mu)} + 1}$ Fermi-Dirac

Fermi energy for non-interacting particles: $E_F = \frac{\hbar^2}{2m} \left(\frac{6\pi^2 n}{g} \right)^{2/3}$
d.o.f. = 1
spin degeneracy (2s+1)

Bose condensation temperature: $T_c = \frac{2\pi\hbar^2}{m k} \left[\frac{n}{g \zeta(3/2)} \right]^{2/3}$
zeta function ≈ 2.612

Electrons in solids

$$\vec{j} = -ne v_d$$

↑
mean electron drift velocity

$$v_d = -\frac{e\tau}{m_e} E$$

↑
applied electric field

mean time between collisions

Fermi gas

electron density of states: $g(E) = \frac{V}{2\pi^2} \left(\frac{2m_e}{\hbar^2}\right)^{3/2} E^{1/2}$

$$g(E_F) = \frac{3}{2} \frac{nV}{E_F}$$

$$k_F = \left(3\pi^2 \frac{N}{V}\right)^{1/3} \quad v_F = \frac{\hbar k_F}{m_e}$$

$$E_F = \frac{\hbar^2 k_F^2}{2m_e} \quad T_F = \frac{E_F}{k_B}$$

$$\frac{C_V}{\text{electron}} = \frac{\pi^2}{3} g(E_F) k_B^2 T = \frac{\pi^2 k_B^2}{2E_F} T$$

$$U_0 = \frac{3}{5} nV E_F$$

$$M = \frac{3n}{2E_F} \mu_B^2 H$$

E & M :

Electric dipole

$$E_r = \frac{p \cos \theta}{2\pi \epsilon_0 r^3}$$

$$E_\theta = \frac{p \sin \theta}{4\pi \epsilon_0 r^3}$$

Image charges

real charge, +q, at a distance .

b from a conducting plane

b from a conducting sphere (radius a)

b from a plane dielectric boundary:

seen from free space

dielectric

image point

-b

a²/b

-b

b

image charge

-q

-q a/b

-q (εr-1)/(εr+1)

2q/εr+1

electric dipole moment, $\vec{p} = q \vec{a}$ or $\vec{p} = \int \vec{r}' \rho d\tau'$

$$\Phi = \frac{\vec{p} \cdot \vec{r}}{4\pi\epsilon_0 r^3}$$

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P}, \epsilon_r = 1 + \chi_E, \vec{D} = \epsilon \vec{E}$$

magnetic dipole moment, $d\vec{m} = \vec{I} ds$, $\vec{m} = \frac{1}{2} \int \vec{r}' \times \vec{J} d\tau'$

$$\vec{B} = \mu_0 (\vec{H} + \vec{M}), \vec{B} = \mu \vec{H}, \mu_r = 1 + \chi_H$$

Force

on a current-carrying element in a magnetic field

$$d\vec{F} = I d\vec{l} \times \vec{B}$$

on a charge

$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$$

on an electric dipole

$$\vec{F} = (\vec{p} \cdot \nabla) \vec{E}$$

on a magnetic dipole

$$\vec{F} = (\vec{m} \cdot \nabla) \vec{B}$$

Torque on an electric dipole

$$\vec{\tau} = \vec{p} \times \vec{E}$$

magnetic "

$$\vec{\tau} = \vec{m} \times \vec{B}$$

on current loops

$$\vec{\tau} = I \oint \vec{r} \times (d\vec{l} \times \vec{B})$$

energy in an electric field $U_{el} = \vec{p} \cdot \vec{E}$
magnetic " $U_{mag} = \vec{m} \cdot \vec{B}$

Hamiltonian of a charged particle in an EM field $H = \frac{|\vec{p}_m - q\vec{A}|^2}{2m} + q\phi$

LCR circuits

$$I = \frac{dQ}{dt}, V = IR, \vec{J} = \sigma \vec{E}, \rho = \frac{1}{\sigma} = \frac{RA}{l}, C = \frac{Q}{V}, I = C \frac{dV}{dt}$$

$$V = -L \frac{dI}{dt}, L_{12} = \frac{\Phi_1}{I_2}$$

energy stored in capacitor $U = \frac{1}{2} CV^2 = \frac{1}{2} QV = \frac{1}{2} \frac{Q^2}{C}$

inductor $U = \frac{1}{2} LI^2 = \frac{1}{2} \frac{\Phi}{I} I = \frac{1}{2} \frac{\Phi^2}{L}$

Power dissipated in a resistor, $W = IV = I^2 R = \frac{V^2}{R}$

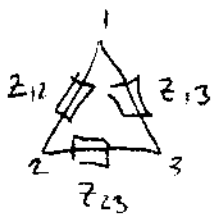
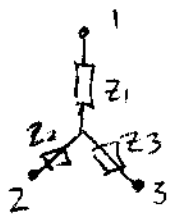
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Impedance of capacitance $Z_C = -\frac{i}{\omega C}$

inductance $Z_L = i\omega L$

Kirchhoff's laws: Current law $\sum_{\text{node}} I_i = 0$ Voltage law $\sum_{\text{loop}} V_i = 0$

Star-delta transformation



$$\leftarrow Z_{ij} = Z_i Z_j \left(\frac{1}{Z_i} + \frac{1}{Z_j} + \frac{1}{Z_k} \right)$$

$$Z_i = \frac{Z_{ij} Z_{ik}}{Z_{ij} + Z_{ik} + Z_{jk}}$$