

$$V(x) = \frac{1}{2} m \omega^2 (x - \epsilon(t))^2 ; \quad \epsilon(t) = \epsilon e^{-t/\tau} \quad \epsilon \ll 1$$

At $t = -\infty$ the particle is in the ground state ($n=0$)

what are the possible states the particle can be in

at $t = \infty$. Work to lowest order of ϵ .

$$V(x) = \frac{1}{2} m \omega^2 (x^2 - 2x\epsilon(t) + \epsilon^2(t)) = \frac{1}{2} m \omega^2 x^2 - m \omega^2 x \epsilon(t) + \frac{1}{2} m \omega^2 \epsilon^2(t)$$

disregard

$$c(t) = \frac{-i}{\hbar} \int_{-\infty}^t \langle H' \rangle e^{i\omega_0 t} ; \quad \omega_0 = \frac{E_1 - E_0}{\hbar}$$

Now $H' = -m\omega^2 x \epsilon e^{-t/\tau}$ and $\langle H' \rangle = \langle n | H' | 0 \rangle$

$$\Rightarrow -m\omega^2 \epsilon e^{-t/\tau} \langle n | x | 0 \rangle ; \quad \text{but } x = \sqrt{\frac{\hbar}{2m\omega}} (a + a^\dagger)$$

$$\Rightarrow -m\omega^2 \epsilon e^{-t/\tau} \left(\sqrt{\frac{\hbar}{2m\omega}} [\langle n | a | 0 \rangle + \langle n | a^\dagger | 0 \rangle] \right)$$

$$= \sqrt{\hbar\omega} \quad \text{for } n=1$$

$$= 0 \quad \text{for } n \neq 1$$

$$= -m\omega^2 \epsilon \sqrt{\frac{\hbar}{2m\omega}} e^{-t/\tau} \quad \text{and } \omega_0 = \frac{E_1 - E_0}{\hbar} = \frac{(1+1/2)\hbar\omega - (0+1/2)\hbar\omega}{\hbar} = \frac{3}{2}\omega - \frac{1}{2}\omega = \omega$$

$$\text{so } c_{0 \rightarrow 1}(t) = \frac{-i}{\hbar} \left(-m\omega^2 \epsilon \sqrt{\frac{\hbar}{2m\omega}} \right) \int_{-\infty}^t e^{-t/\tau} e^{i\omega t} dt$$

we need to complete the square:

$$\int_{-\infty}^t e^{-(t/\tau - i\omega t)} dt \quad y \equiv \sqrt{a} x + \frac{b}{2\sqrt{a}} \Rightarrow y^2 = ax^2 + bx + \frac{b^2}{4a}$$

$$\text{so } a = \frac{1}{\tau^2} ; \quad b = -i\omega \Rightarrow y = \frac{1}{\tau} x - \frac{i\omega}{2} \quad dy = \frac{1}{\tau} dx$$

$$\text{so } \int_{-\infty}^{\infty} e^{-(\alpha^2 x^2 - i\omega x)} dx = \frac{1}{\tau^{-1}} \int_{-\infty}^{\infty} e^{-(y^2 - \frac{b^2}{4a})} dy = \frac{e^{b^2/4a}}{\tau^{-1}} \underbrace{\int_{-\infty}^{\infty} e^{-y^2} dy}_{\sqrt{\pi}} = \frac{\sqrt{\pi}}{\tau^{-1}} e^{b^2/4a}$$

$$2 \int_0^{\infty} e^{-y^2} dy = \sqrt{\pi}$$

hence $c_{0 \rightarrow 1}(\omega) = \frac{i m \omega^3 \epsilon \sqrt{\pi}}{\hbar^2 \tau^{-1} \sqrt{2 m \omega}} e^{b^2/4a}$

and the probability is: $|c_{0 \rightarrow 1}|^2 = \frac{m^2 \omega^6 \epsilon^2 \pi}{\hbar^2 \tau^{-2} 2 m \omega} e^{b^2/2a}$

$$= \frac{m \omega^3 \epsilon^2 \pi}{2 \hbar \tau^{-2}} e^{(-i\omega)^2/2(\frac{1}{\tau^2})} = \frac{m \omega^3 \epsilon^2 \pi}{2 \hbar \tau^{-2}} e^{-\frac{\omega^2 \tau^2}{2}} = \frac{m \omega^3 \epsilon^2 \tau^2 \pi}{2 \hbar} e^{-\frac{\omega^2 \tau^2}{2}}$$