

A nonrelativistic particle of mass m is bound in a central potential $V(\vec{r}) = V(r)$ that is finite at the origin.

- (a) Show that its wave function in spherical coordinates can be written in the factorized form $\psi(\vec{r}) = R(r)Y(\theta, \phi)$. Show that the radial part T_r can be expressed in terms of the radial part \vec{p}_r of the linear momentum operator \vec{p} , and that
- $$\vec{p}_r \equiv \frac{1}{2}(\hat{e}_r \cdot \vec{p} + \vec{p} \cdot \hat{e}_r) = \frac{\hbar}{i} \frac{1}{r} \frac{d}{dr} r$$

→ (see Griffiths pgs 121-133)

In spherical coordinates, the Schrödinger equation is

$$-\frac{\hbar^2}{2m} \left[\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{d\psi}{dr} \right) + \frac{1}{r^2 \sin\theta} \frac{d}{d\theta} \left(\sin\theta \frac{d\psi}{d\theta} \right) + \frac{1}{r^2 \sin^2\theta} \left(\frac{d^2\psi}{d\phi^2} \right) \right] + V\psi = E\psi$$

try solution $\psi(r, \theta, \phi) = R(r)Y(\theta, \phi)$

$$-\frac{\hbar^2}{2m} \left[\frac{Y}{r^2} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \frac{R}{r^2 \sin\theta} \frac{d}{d\theta} \left(\sin\theta \frac{dY}{d\theta} \right) + \frac{1}{r^2 \sin^2\theta} \left(\frac{d^2Y}{d\phi^2} \right) \right] + VRY = ERY$$

Divide by RY and multiply by $-\frac{2mr^2}{\hbar^2}$ to yield:

$$\underbrace{\left[\frac{1}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) - \frac{2mr^2}{\hbar^2} (V(r) - E) \right]}_{=l(l+1)} + \underbrace{\frac{1}{Y} \left[\frac{1}{\sin\theta} \frac{d}{d\theta} \left(\sin\theta \frac{dY}{d\theta} \right) + \frac{1}{\sin^2\theta} \frac{d^2Y}{d\phi^2} \right]}_{=-l(l+1)} = 0$$

each term in brackets must be constant, let separation constant be $l(l+1)$
now verify that $\frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) \propto \vec{p}_r^2 R$

$$\vec{p}_r^2 R = -\hbar^2 \left(\frac{1}{r} \frac{d}{dr} r \right) \left(\frac{1}{r} \frac{d}{dr} (rR) \right) = -\hbar^2 \left(\frac{1}{r} \frac{d}{dr} r \right) \frac{1}{r} \left(R + r \frac{dR}{dr} \right)$$

$$= -\hbar^2 \frac{1}{r} \frac{d}{dr} \left(R + r \frac{dR}{dr} \right) = -\hbar^2 \left(\frac{dR}{dr} + \frac{dR}{dr} + r \frac{d^2R}{dr^2} \right)$$

$$= -\hbar^2 \left(\frac{2}{r} \frac{dR}{dr} + \frac{d^2R}{dr^2} \right) = -\hbar^2 \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right)$$

thus, one can write $T_r = \frac{\vec{p}_r^2}{2mr^2}$

b) Obtain the normalized angular function $Y(\theta, \phi)$ for S-states. Show that the S-state radial function $R(r)$ could have solutions near the origin of the forms $R_1(r) = A$ and $R_2(r) = B/r$ where A, B are real constants. Show that both solutions are square-integrable near the origin.

→ for S-states, $l=0$, and the angular equation becomes

$$\sin \theta \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial Y}{\partial \theta} \right) + \frac{\partial^2 Y}{\partial \phi^2} = 0$$

→ Y must be constant, normalization yields

$$Y = \frac{1}{\sqrt{4\pi}}$$

for $R_1(r) = A$:

$$\frac{1}{A} \frac{d}{dr} \left(r^2 \frac{dA}{dr} \right) - \frac{2mr^2}{\hbar^2} (V(r) - E) = 0$$

near origin: $\int_0^{R_0} A^2 r^2 dr = \frac{A^2 R_0^3}{3}$

where R_0 is small

$$-\frac{2mr^2}{\hbar^2} [V(r) - E] = 0$$

for $R_2(r) = B/r$:

$$\frac{1}{B} \frac{d}{dr} \left(r^2 \frac{d}{dr} \left(\frac{B}{r} \right) \right) - \frac{2mr^2}{\hbar^2} (V(r) - E) = 0$$

near origin: $\int_0^{R_0} \frac{B^2}{r^2} r^2 dr = B^2 R_0$

$$-\frac{2mr^2}{\hbar^2} (V(r) - E) = 0$$

c) Show that $R_2(r)$ is not admissible as a solution for the wave function because the radial momentum \hat{p}_r is not Hermitian for this type of wave function.

$$\int_0^{\infty} \frac{B}{r} \left(\hat{p}_r \left(\frac{B}{r} \right) \right) r^2 dr = \int_0^{\infty} \frac{B}{r} \frac{\hbar}{i} \frac{1}{r} \left(\frac{d}{dr} \left(r \frac{B}{r} \right) \right) r^2 dr = 0$$

thus, \hat{p}_r is not Hermitian