

Here are the equations of states for an ideal gas:

$$PV = Nk_B T; \quad U = \frac{3}{2} Nk_B T,$$

where P is the pressure of the gas, N is the number of particles, T is the temperature, U is the internal energy, and k_B is Boltzmann's constant. Derive the Helmholtz free energy, $F(N, V, T)$, of the ideal gas. From this quantity, find the entropy as a function of temperature, number of particles, and volume.

See Reif. pgs. 238-245 and p. 50 for his definition

$$Z_1 = \frac{1}{h_0^3} \int_V \int_{-\infty}^{\infty} e^{-\beta E} d^3r d^3p \quad \text{now } E = \frac{\vec{p}^2}{2m} \text{ for an ideal gas}$$

$$= \frac{1}{h_0^3} \underbrace{\int_V d^3r}_{=V} \int_{-\infty}^{\infty} e^{-\frac{\beta}{2m} \vec{p}^2} d^3p = \frac{V}{h_0^3} \int_{-\infty}^{\infty} e^{-\frac{\beta}{2m} \vec{p}^2} d^3p$$

$$\text{now } \vec{p}^2 = p_x^2 + p_y^2 + p_z^2, \text{ so } \int_{-\infty}^{\infty} e^{-\frac{\beta}{2m} \vec{p}^2} d^3p = \int_{-\infty}^{\infty} e^{-\frac{\beta}{2m} p_x^2} dp_x \int_{-\infty}^{\infty} e^{-\frac{\beta}{2m} p_y^2} dp_y \int_{-\infty}^{\infty} e^{-\frac{\beta}{2m} p_z^2} dp_z$$

we are dealing with an integral of the form $\int_0^{\infty} e^{-ax^2} dx = \frac{1}{2} \sqrt{\frac{\pi}{a}}$

$$\text{so } \int_{-\infty}^{\infty} e^{-\frac{\beta}{2m} p_x^2} dp_x = 2 \int_0^{\infty} e^{-\frac{\beta}{2m} p_x^2} dp_x = 2 \cdot \frac{1}{2} \sqrt{\frac{\pi}{\beta/2m}} = \sqrt{\frac{2\pi m}{\beta}} = \sqrt{2\pi m k_B T}$$

$$\text{hence } \int_{-\infty}^{\infty} e^{-\frac{\beta}{2m} \vec{p}^2} d^3p = (2\pi m k_B T)^{3/2}$$

$$\text{and then } Z_1 = \frac{V}{h_0^3} (2\pi m k_B T)^{3/2} = \frac{V}{\lambda_{th}^3}; \quad \lambda_{th} = \sqrt{\frac{h_0^2}{2\pi m k_B T}}$$

thermal de Broglie (similar to it actually)
as $h_0 \rightarrow h$
↑
classical → quantum

$$\text{So } Z = \frac{z_1^N}{N!}$$

Stirling's Approx.

$$\ln N! = N \ln N - N$$

$$\text{now } E = -kT \ln Z = -kT \ln \frac{z_1^N}{N!} = -kT \left[N \ln z_1 - \ln N! \right]$$

$$\downarrow$$

$$\uparrow \frac{V}{2h^3}$$

$$= -NkT \left[\ln \frac{V}{2h^3} - \ln N + 1 \right] = -NkT \left[\ln V - \ln N - 3 \ln 2 + 1 \right]$$

$$= -NkT \left[\ln \frac{V}{N} - \frac{3}{2} \ln \frac{h_0^3}{2\pi m k T} + 1 \right]$$

$$\text{as for } S: S = -\left(\frac{\partial F}{\partial T}\right)_{V,N} = Nk \frac{\partial}{\partial T} T \left[\ln \frac{V}{N} - \frac{3}{2} \ln \frac{h_0^3}{2\pi m k T} - 1 \right]$$

$$= Nk \left(\left[\ln \frac{V}{N} - \frac{3}{2} \ln \frac{h_0^3}{2\pi m k T} + 1 \right] + T \left[-\frac{3}{2} \left(\frac{1}{\frac{h_0^3}{2\pi m k T}} \right) \cdot \left(\frac{-h_0^3}{2\pi m k T^2} \right) \right] \right)$$

$$= Nk \left(\left[\ln \frac{V}{N} - \frac{3}{2} \ln \frac{h_0^3}{2\pi m k T} - 1 \right] + T \left[\frac{+3}{2} \frac{1}{T} \right] \right)$$

$$= Nk \left[\ln \frac{V}{N} - \frac{3}{2} \ln \frac{h_0^3}{2\pi m k T} + \frac{5}{2} \right]$$