

# Spring 1998 #1 (p 1 of 2)

① We want to solve  $\nabla \cdot \vec{D} = \epsilon_0 \nabla \cdot \vec{E} = \rho$ ,  $z > 0$   
 $\nabla \cdot \vec{D} = \epsilon \nabla \cdot \vec{E} = 0$ ,  $z < 0$   
 $\nabla \times \vec{E} = 0$

and due to the boundary conditions we require

$$\lim_{z \rightarrow 0^-} \begin{Bmatrix} \epsilon E_z \\ E_p \end{Bmatrix} = \lim_{z \rightarrow 0^+} \begin{Bmatrix} \epsilon_0 E_z \\ E_p \end{Bmatrix} \quad (\text{clearly no } E_\phi)$$

Since  $\nabla \times \vec{E} = 0$  write down  $\Phi$ . Try obvious image charges:

②  $z > 0$ : image  $Q'$  at  $z = -d$ .  $\Phi = \frac{1}{4\pi\epsilon_0} \left( \frac{Q}{R_1} + \frac{Q'}{R_2} \right)$

with  $R_1 = \sqrt{(z-d)^2 + \rho^2}$   $R_2 = \sqrt{(z+d)^2 + \rho^2}$

then  $\vec{E} = -\nabla \Phi = -\left( \hat{\rho} \frac{\partial}{\partial \rho} + \frac{1}{\rho} \hat{\phi} \frac{\partial}{\partial \phi} + \hat{z} \frac{\partial}{\partial z} \right) \Phi = \frac{1}{4\pi\epsilon_0} \left( \hat{\rho} \left( \frac{\partial}{\partial \rho} \frac{Q}{R_1} + \frac{\partial}{\partial \rho} \frac{Q'}{R_2} \right) + \left( \hat{z} \left( \frac{\partial}{\partial z} \frac{Q}{R_1} + \frac{\partial}{\partial z} \frac{Q'}{R_2} \right) \right) \right)$   
 $= \frac{1}{4\pi\epsilon_0} \left( \hat{\rho} \left( Q \frac{\rho}{((z-d)^2 + \rho^2)^{3/2}} + Q' \frac{\rho}{((z+d)^2 + \rho^2)^{3/2}} \right) + \hat{z} \left( Q \frac{z-d}{((z-d)^2 + \rho^2)^{3/2}} + Q' \frac{z+d}{((z+d)^2 + \rho^2)^{3/2}} \right) \right)$

③  $z < 0$ : image  $Q''$  at  $z = +d$ :

$$\Phi = \frac{1}{4\pi\epsilon} \frac{Q''}{R_1} \Rightarrow \vec{E} = +\frac{1}{4\pi\epsilon} \left( \hat{\rho} Q'' \frac{\rho}{((z-d)^2 + \rho^2)^{3/2}} + \hat{z} Q'' \frac{z-d}{((z-d)^2 + \rho^2)^{3/2}} \right)$$

at  $z=0$ , require  $\epsilon_0 E_z(z>0) = \epsilon E_z(z<0) \Rightarrow \frac{1 \cdot \epsilon_0}{4\pi\epsilon_0} \left( Q \frac{-d}{(d^2 + \rho^2)^{3/2}} + Q' \frac{d}{(d^2 + \rho^2)^{3/2}} \right) = \frac{1 \cdot \epsilon}{4\pi\epsilon} \left( Q'' \frac{-d}{(d^2 + \rho^2)^{3/2}} \right)$

$$\therefore \boxed{+Q - Q' = +Q''}$$

also require  $E_p = E_p$ :  $\frac{1}{4\pi\epsilon_0} \left( \frac{Q\rho}{(d^2 + \rho^2)^{3/2}} + \frac{Q'\rho}{(d^2 + \rho^2)^{3/2}} \right) = \frac{1}{4\pi\epsilon} \left( \frac{Q''\rho}{(d^2 + \rho^2)^{3/2}} \right) \Rightarrow \epsilon(Q + Q') = \epsilon_0 Q''$

# Spring 1998 #1 (p 2 of 2)

we care about the energy, i.e. the force, i.e. the field on  $Q$ , which is in  $(I)$ , so

$$Q - Q' = Q'' = \frac{\epsilon}{\epsilon_0} Q + \frac{\epsilon}{\epsilon_0} Q' \Rightarrow Q \left(1 - \frac{\epsilon}{\epsilon_0}\right) = Q' \left(\frac{\epsilon}{\epsilon_0} + 1\right) \Rightarrow Q' = \frac{\epsilon_0 - \epsilon}{\epsilon_0 + \epsilon} Q$$

could solve for the field here, but force between 2 point charges is just:

$$F = \frac{+1}{4\pi\epsilon_0} \frac{QQ'}{r^2} = \frac{-Q^2}{16\pi\epsilon_0 d^2} \left(\frac{\epsilon_0 - \epsilon}{\epsilon_0 + \epsilon}\right) \quad (\text{negative because attractive})$$

and the work to assemble this is just

$$U = -\left(\frac{\epsilon_0 - \epsilon}{\epsilon_0 + \epsilon}\right) \frac{Q^2}{4\pi\epsilon_0} \int_{\infty}^{2d} \frac{1}{r^2} dr = -\left(\frac{\epsilon_0 - \epsilon}{\epsilon_0 + \epsilon}\right) \frac{Q^2}{4\pi\epsilon_0} \left(-\frac{1}{2d}\right)$$

For fun, the long way:  $p=0$  axis component of  $E$ :

$$E_z = \frac{+Q}{4\pi\epsilon_0} \left( \frac{z-d}{(z-d)^3} + \frac{(\epsilon_0 - \epsilon)}{(\epsilon_0 + \epsilon)} \frac{z+d}{(z+d)^3} \right)$$

so that at  $z=d$ , the force is  $F = QE = \frac{-Q^2}{4\pi\epsilon_0} \left(\frac{\epsilon_0 - \epsilon}{\epsilon_0 + \epsilon}\right) \frac{1}{4d^2} = \frac{-Q^2}{16\pi\epsilon_0 d^2} \left(\frac{\epsilon_0 - \epsilon}{\epsilon_0 + \epsilon}\right)$

and  $U = \frac{-Q^2}{4\pi\epsilon_0} \int_{\infty}^d \frac{1}{(z+d)^3} \left(\frac{\epsilon_0 - \epsilon}{\epsilon_0 + \epsilon}\right) dz = \frac{-Q^2}{4\pi\epsilon_0} \left(-\frac{1}{2z}\right) \Big|_{\infty}^d \left(\frac{\epsilon_0 - \epsilon}{\epsilon_0 + \epsilon}\right)$  (knowing self-energy)

or just use  $U = Q\Phi = \frac{Q^2}{4\pi\epsilon_0} \left(\frac{\epsilon_0 - \epsilon}{\epsilon_0 + \epsilon}\right) \frac{1}{R_1} \Big|_{\infty}^d = \frac{Q^2}{4\pi\epsilon_0} \left(\frac{\epsilon_0 - \epsilon}{\epsilon_0 + \epsilon}\right) \frac{1}{2d}$