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Consider the motion of a charge in an electric field  $\vec{E}$ , using the correct equations of special relativity. Let  $\vec{E} = E_0 \hat{x}$ , with  $E_0$  constant. Consider a particle with initial conditions

$$\begin{aligned} x=y=z=v_x=v_z=0 & \quad \text{at } t=0 \\ v_y \neq 0 & \quad \text{at } t=0 \end{aligned}$$

a) Solve for  $v_x(t)$ ,  $v_y(t)$ , and  $v_z(t)$ . Also find simplified expressions for  $v_x$  and  $v_y$  when  $t \rightarrow \infty$ .

From Newton's 2nd law, we have

$$\frac{d\vec{p}}{dt} = \vec{F} = q\vec{E} + q\vec{v} \times \vec{B} = q\vec{E} = qE_0 \hat{x} \quad (1)$$

where  $p_\mu$  is the 4-momentum defined a-like so

$$p_\mu = p_0 + \vec{p} = (E, p_x, p_y, p_z) \quad \text{note: } c=1 \text{ here} \quad (2)$$

From eq (1), we have

$$\begin{aligned} \frac{dp_x}{dt} &= qE_0, \quad \frac{dp_y}{dt} = 0 = \frac{dp_z}{dt} \\ \Rightarrow p_x &= qE_0 t, \quad p_y = c_1, \quad p_z = c_2 \end{aligned}$$

From initial conditions, these values become

$$p_x = qE_0 t, \quad p_y = p_{yi}, \quad p_z = 0$$

From Abers eq (12.76), we have, in natural units,

$$\vec{v} = \frac{\vec{p}}{E} = \frac{\vec{p}}{\sqrt{p^2 + m^2}} \quad (3)$$

But, what is  $p^2$ ?

$$p^2 = \vec{p} \cdot \vec{p} = p_x p_x + p_y p_y + p_z p_z = (qE_0 t)^2 + p_{yi}^2 + 0 \quad (4)$$

substituting this result into eq. (3) yields

$$\vec{v} = \frac{\vec{p}}{\sqrt{(qE_0 t)^2 + p_{yi}^2 + m^2}}$$

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note:  $\sqrt{p_{yi}^2 + m^2}$  is the initial energy,  $\epsilon_0$

So, we have

$$\vec{v} = \frac{\vec{p}}{\sqrt{\epsilon_0^2 + (qEt)^2}}$$

Now, we can immediately write down  $v_x$ ,  $v_y$ , and  $v_z$

$$v_x = \frac{p_x}{\sqrt{\epsilon_0^2 + (qEt)^2}} = \frac{qEt}{\sqrt{\epsilon_0^2 + (qEt)^2}} = \frac{1}{\sqrt{1 + \left(\frac{\epsilon_0}{qEt}\right)^2}}$$

$$\Rightarrow \lim_{t \rightarrow \infty} v_x = 1$$

$$v_y = \frac{p_{yi}}{\sqrt{\epsilon_0^2 + (qEt)^2}} = \frac{p_{yi}}{qEt} \frac{1}{\sqrt{1 + \left(\frac{\epsilon_0}{qEt}\right)^2}} = \frac{p_{yi}}{qEt} v_x$$

$$\Rightarrow \lim_{t \rightarrow \infty} v_y = 0$$

$$v_z = 0$$

(b) Find  $x(t)$ ,  $y(t)$ , and find the trajectory  $x(y)$ .

(i)  $x(t)$

From part (a), we have

$$v_x = \frac{dx}{dt} = \frac{qEt}{\sqrt{\epsilon_0^2 + (qEt)^2}} \Rightarrow x(t) = \int_0^t \frac{qEt'}{\sqrt{(qEt')^2 + \epsilon_0^2}} dt'$$

$$\text{let } u = qEt \Rightarrow du = qE dt$$

so, we have

$$x(t) = \frac{1}{qE_0} \int_0^{qEt} \frac{u du}{\sqrt{u^2 + \epsilon_0^2}} = \frac{1}{qE_0} \left[ (u^2 + \epsilon_0^2)^{1/2} \right]_0^{qEt}$$

$$\Rightarrow x(t) = \frac{1}{qE_0} \left[ \sqrt{(qE_0 t)^2 + \epsilon_0^2} - \epsilon_0 \right] \quad (5)$$

note:  $x(0) = \frac{1}{qE_0} [\epsilon_0 - \epsilon_0] = 0$  as required by initial conditions

(ii)  $y(t)$

from part (i), we know

$$v_y = \frac{dy}{dt} = \frac{p_{yi}}{\sqrt{\epsilon_0^2 + (qE_0 t)^2}} \Rightarrow y(t) = \int_0^t \frac{p_{yi} dt'}{\sqrt{\epsilon_0^2 + (qE_0 t')^2}}$$

once again

let  $u = qE_0 t$

so,

$$y(t) = \frac{1}{qE_0} \int_0^{qE_0 t} \frac{p_{yi} du}{\sqrt{\epsilon_0^2 + u^2}} = \frac{p_{yi}}{qE_0} \int_0^{qE_0 t} \frac{du}{\sqrt{\epsilon_0^2 + u^2}}$$

$$= \frac{p_{yi}}{qE_0} \left[ \sinh^{-1} \left( \frac{u}{\epsilon_0} \right) \right]_0^{qE_0 t} = \frac{p_{yi}}{qE_0} \left[ \sinh^{-1} \left( \frac{qE_0 t}{\epsilon_0} \right) - \cancel{\sinh^{-1}(0)} \right]$$

$$\Rightarrow y(t) = \frac{p_{yi}}{qE_0} \sinh^{-1} \left( \frac{qE_0 t}{\epsilon_0} \right) \quad (6)$$

note:  $y(0) = \frac{p_{yi}}{qE_0} \sinh^{-1}(0) = 0$  as required by initial conditions

(iii)  $x(y)$

solve eq. (5) for  $t$ :

$$(qE_0 t)^2 + \epsilon_0^2 = (qE_0 x(t) + \epsilon_0)^2$$

$$\Rightarrow \quad \gamma E_0 t = \sqrt{(\gamma E_0 x(t) + \epsilon_0)^2 - \epsilon_0^2}$$

$$\therefore \quad t = \frac{1}{\gamma E_0} \sqrt{(\gamma E_0 x(t) + \epsilon_0)^2 - \epsilon_0^2} \quad (7)$$

Now, solve eq (6) for t:

$$\frac{\gamma E_0 t}{\epsilon_0} = \sinh \left[ \frac{\gamma E_0 y(t)}{p_{yi}} \right]$$

$$\therefore \quad t = \frac{\epsilon_0}{\gamma E_0} \sinh \left[ \frac{\gamma E_0 y(t)}{p_{yi}} \right] \quad 8$$

Now, set eq (7) equal to eq (8) and solve for  $x(t)$ .

$$\frac{1}{\gamma E_0} \sqrt{(\gamma E_0 x(t) + \epsilon_0)^2 - \epsilon_0^2} = \frac{\epsilon_0}{\gamma E_0} \sinh \left[ \frac{\gamma E_0 y(t)}{p_{yi}} \right]$$

$$\Rightarrow (\gamma E_0 x(t) + \epsilon_0)^2 - \epsilon_0^2 = \epsilon_0^2 \sinh^2 \left[ \frac{\gamma E_0 y(t)}{p_{yi}} \right]$$

$$\begin{aligned} \Rightarrow \quad \gamma E_0 x(t) + \epsilon_0 &= \sqrt{\epsilon_0^2 \left[ \sinh^2 \left( \frac{\gamma E_0 y(t)}{p_{yi}} \right) + 1 \right]} \\ &= \sqrt{\epsilon_0^2 \cosh^2 \left( \frac{\gamma E_0 y(t)}{p_{yi}} \right)} = \epsilon_0 \cosh \left( \frac{\gamma E_0 y(t)}{p_{yi}} \right) \end{aligned}$$

$$\Rightarrow \quad x(t) = \frac{1}{\gamma E_0} \left[ \epsilon_0 \left( \cosh \left( \frac{\gamma E_0 y(t)}{p_{yi}} \right) - 1 \right) \right]$$

$$\therefore \quad x(y) = \frac{\epsilon_0}{\gamma E_0} \left[ \cosh \left( \frac{\gamma E_0 y(t)}{p_{yi}} \right) - 1 \right]$$